

Microensing:

Theory, Practice, Results, Future

Lecture 4

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Outline of Lectures

Lecture 1: Microlensing History and Theory

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- Background
- Motivation/Goals
- Early results
- Evolution of a field
- Basic microlensing theory

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Lecture 1: Microlensing History and Theory

Lecture 2: Beyond the Single lens

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Lecture 2: Beyond the Single lens

- Finite source star
- Limb Darkening
- Blending
- Parallax
- Xallarap

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Lecture 2: Beyond the Single lens

Lecture 3: Planetary Microlensing - I

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Lecture 1: Microlensing History and Theory

Lecture 2: Beyond the Single lens

Lecture 3: Planetary Microlensing - I

- Binary lens microlensing
- Extreme mass ratio microlensing
- Theoretical tools of trade: caustics
- Planetary microlensing regimes
- General rules

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Lecture 1: Microlensing History and Theory

Lecture 2: Beyond the Single lens

Lecture 3: Planetary Microlensing - I

Lecture 4: Planetary Microlensing - II

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Lecture 2: Beyond the Single lens

Lecture 3: Planetary Microlensing - I

Lecture 4: Planetary Microlensing - II

- Capabilities, detection limits
- Detection
- Modelling

Detection techniques

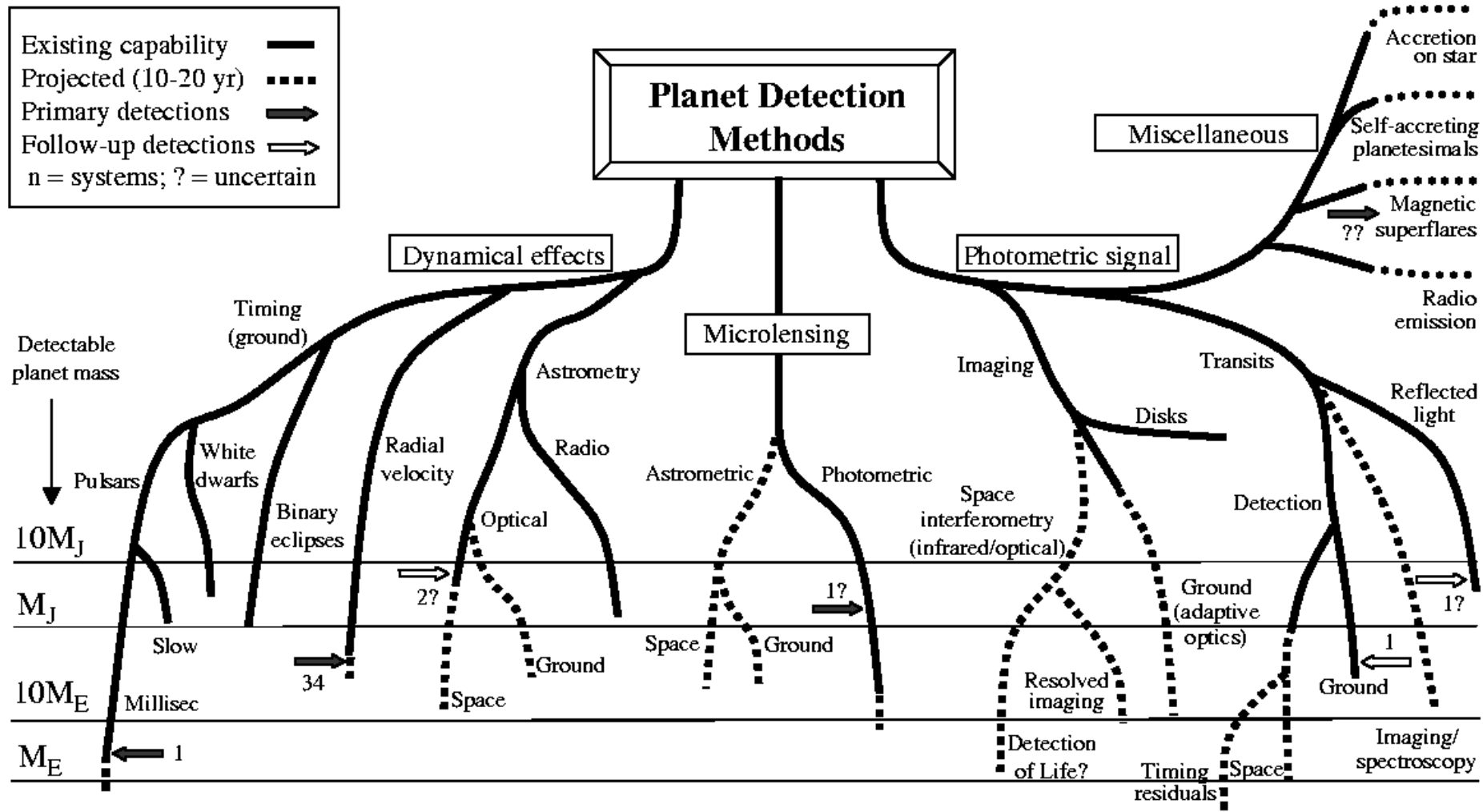
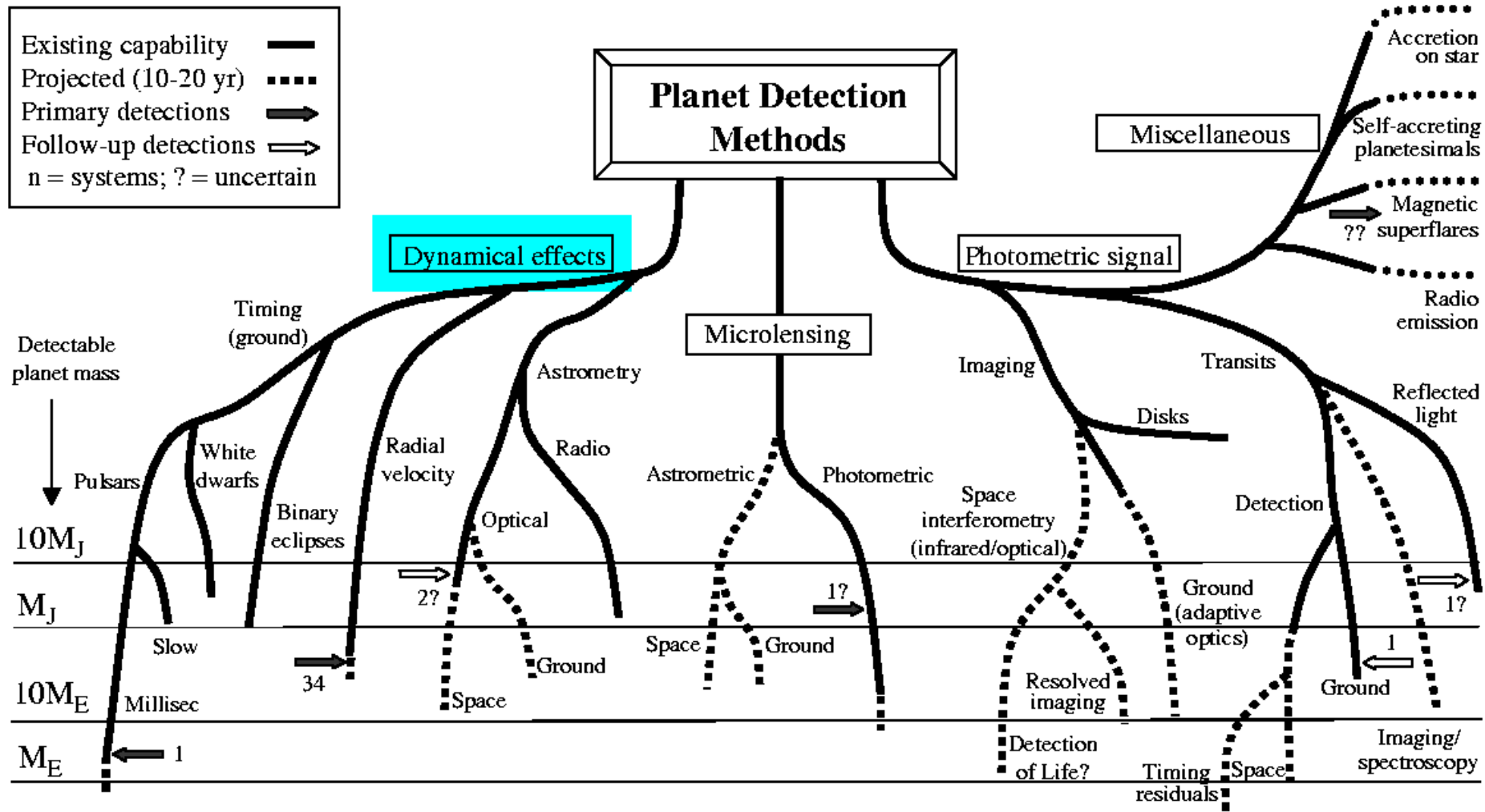


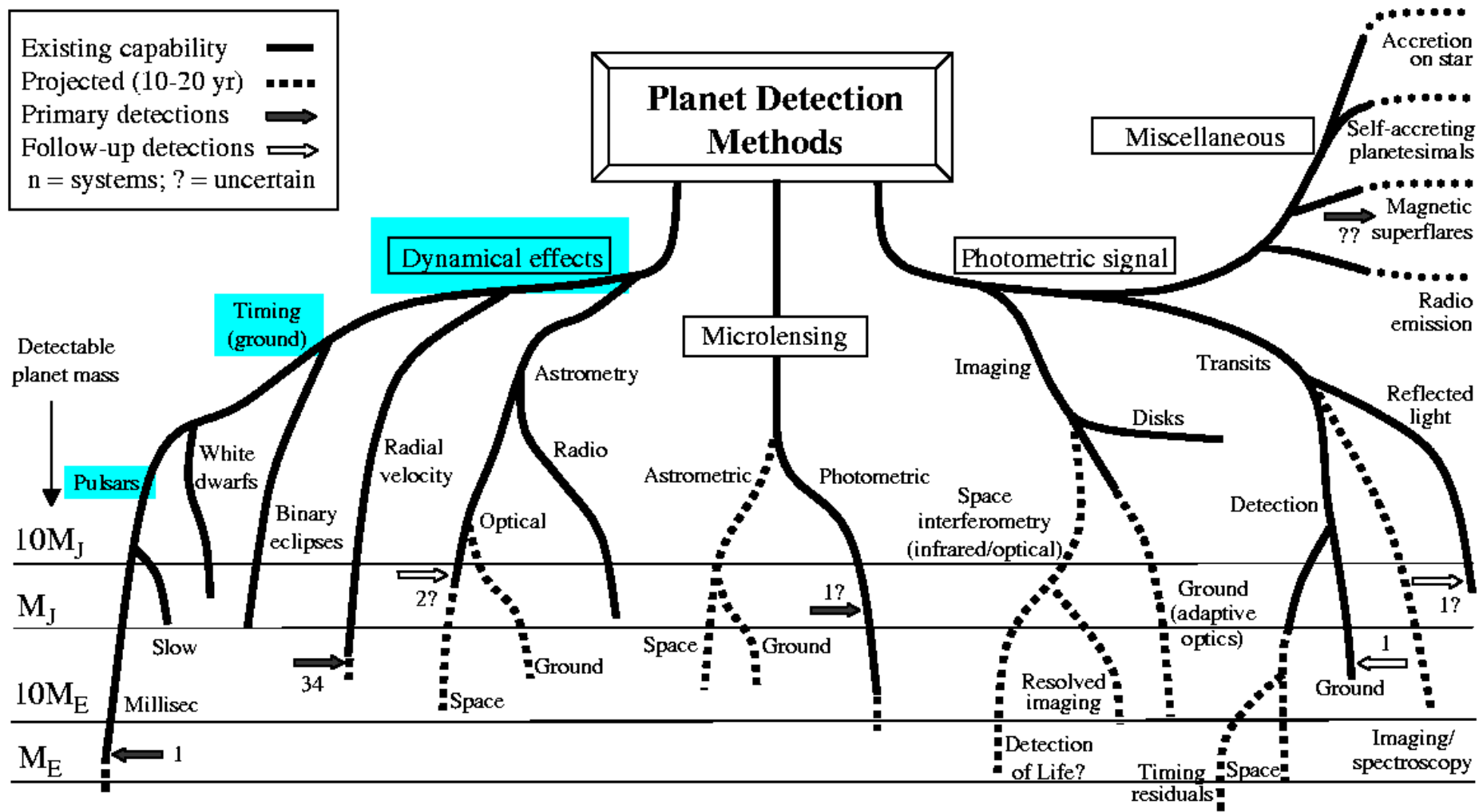
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Detection techniques



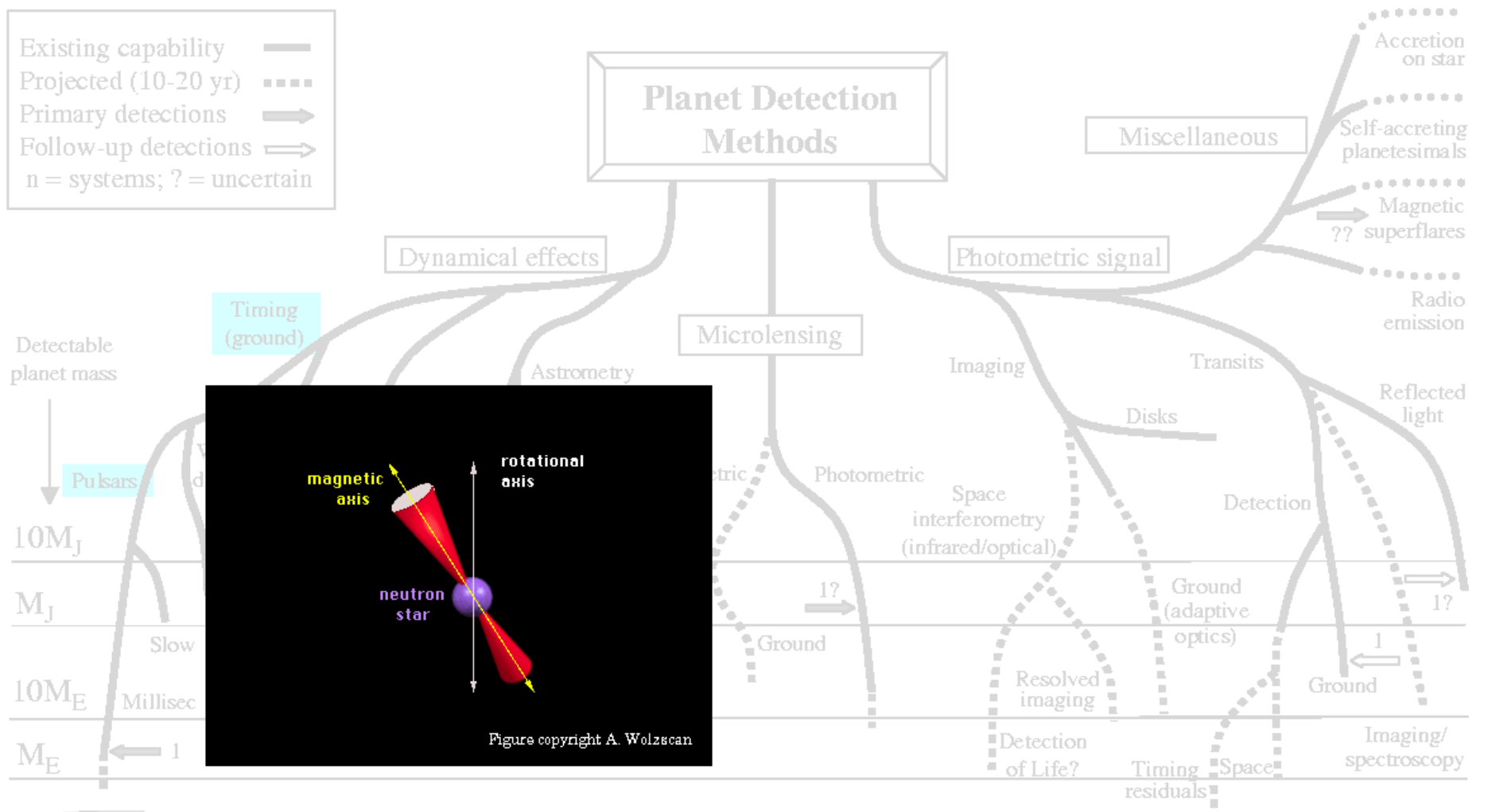
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Existing capability —
 Projected (10-20 yr) ·····
 Primary detections →
 Follow-up detections ⇨
 n = systems; ? = uncertain

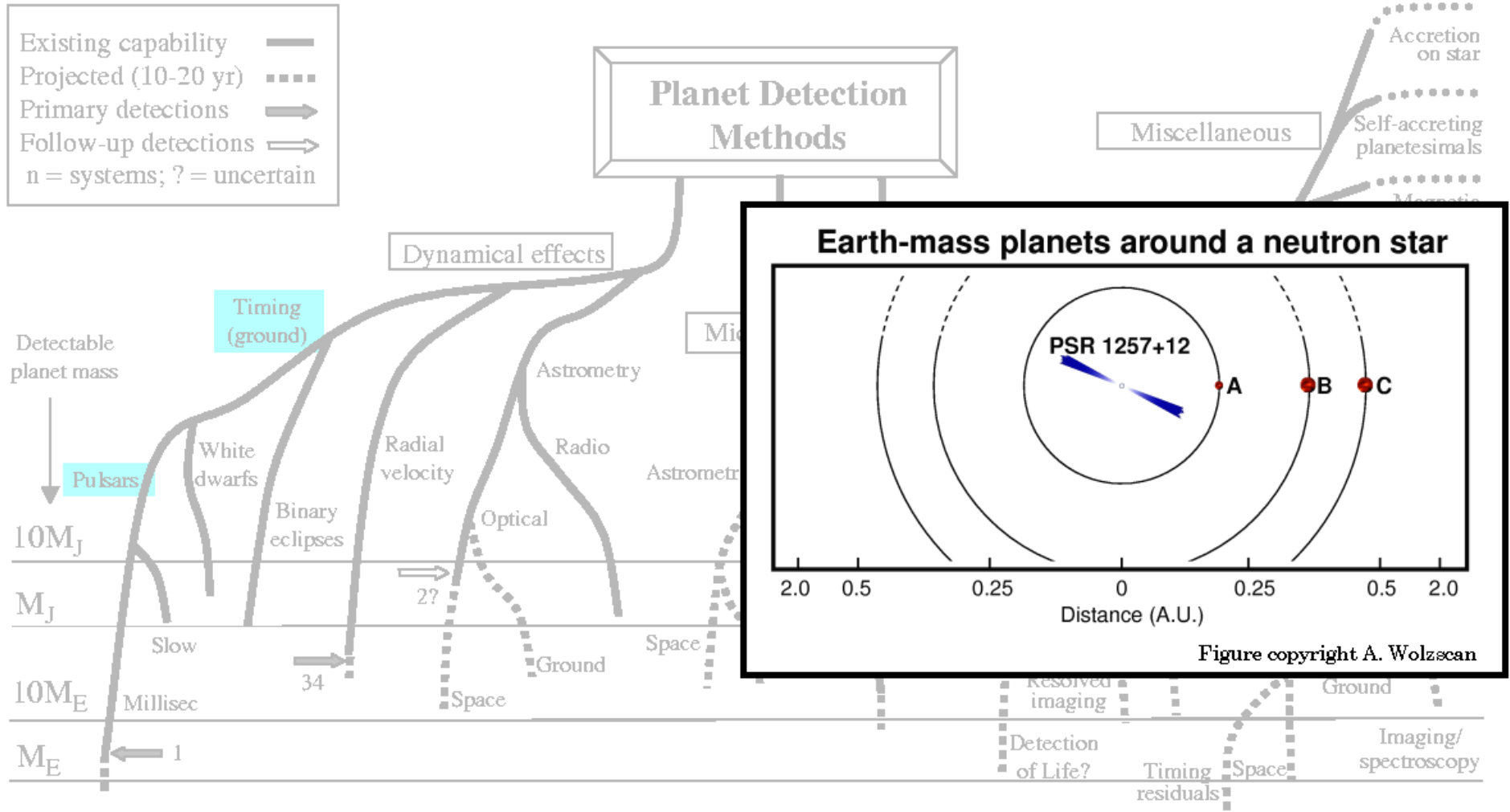


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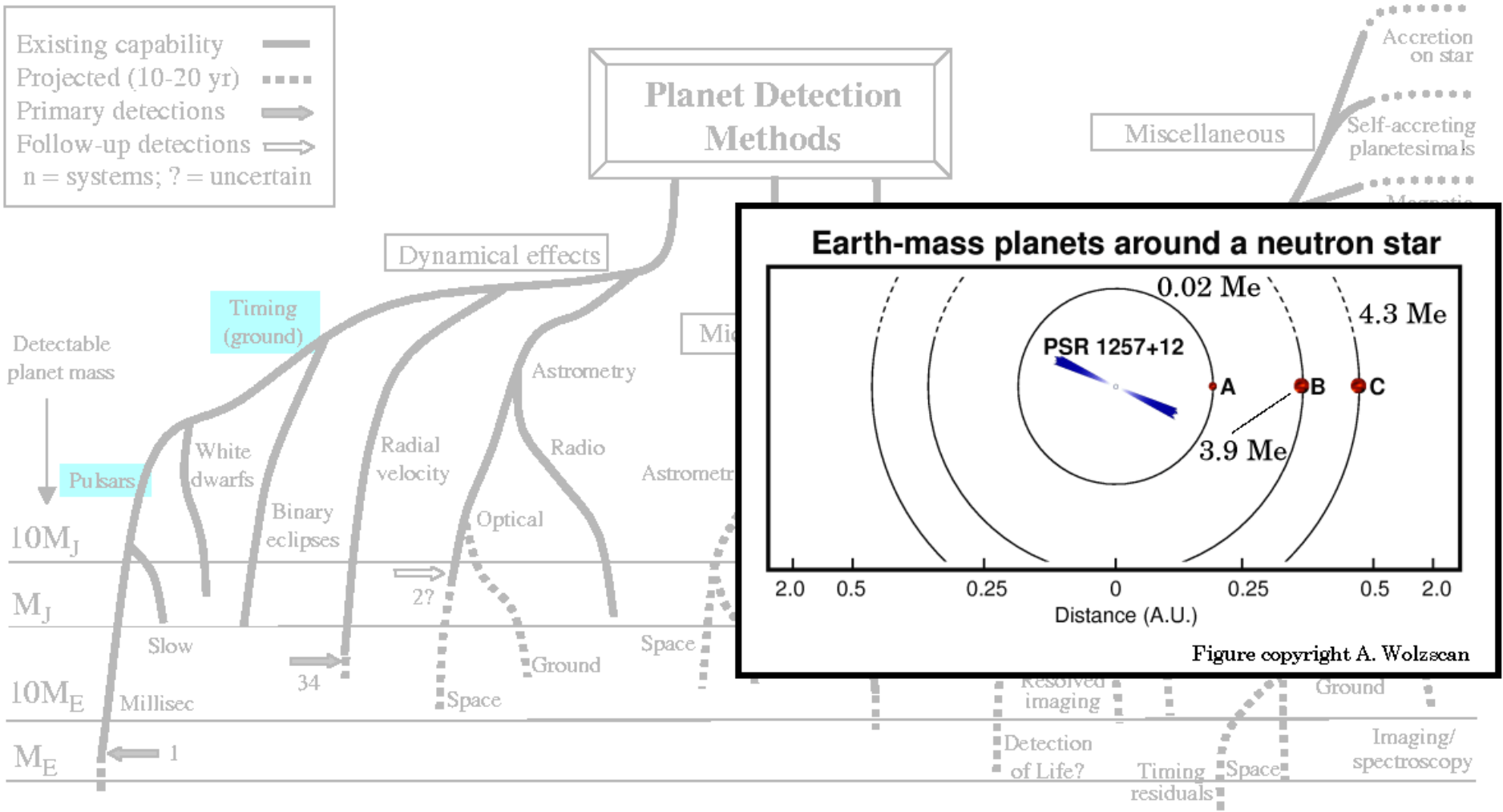
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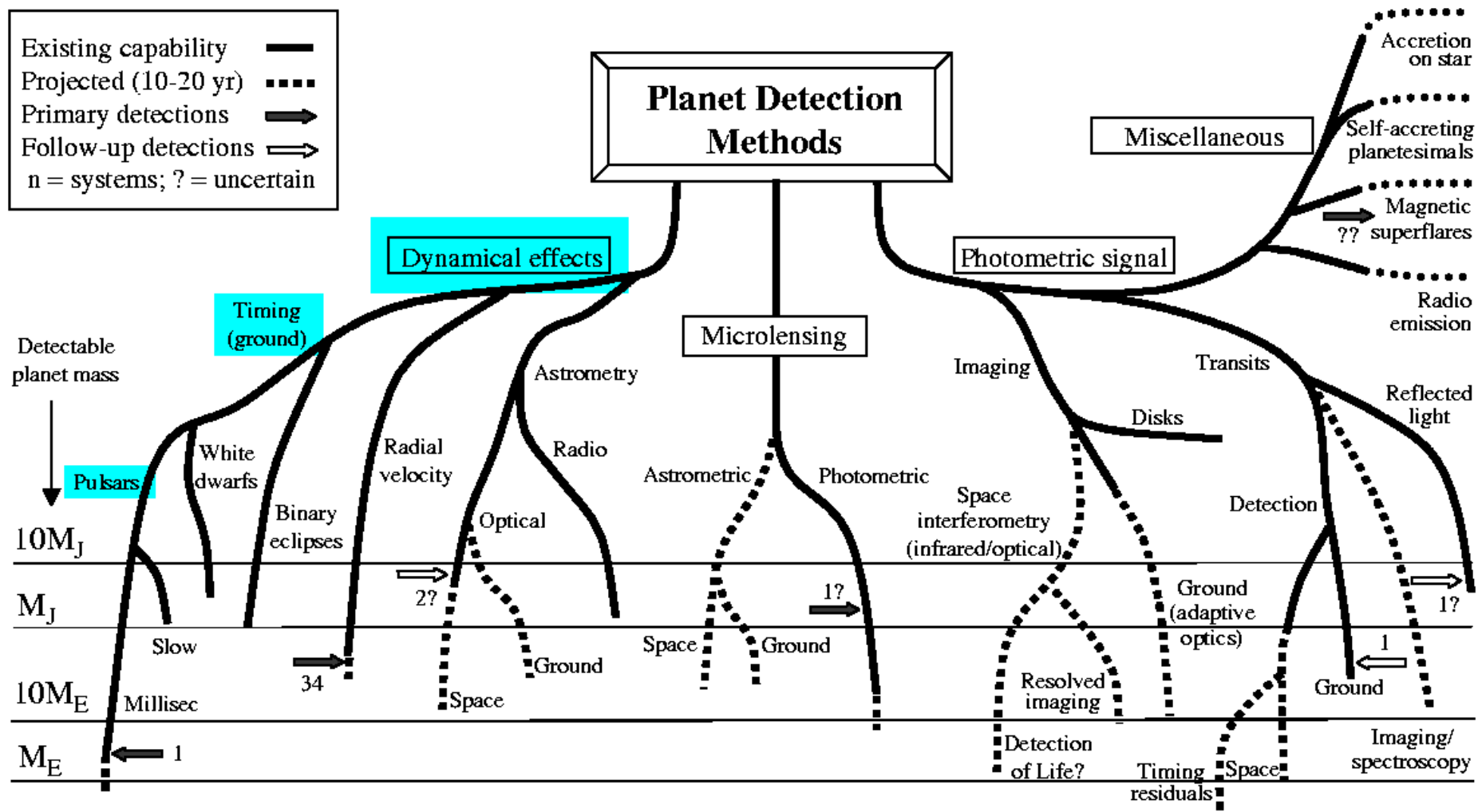


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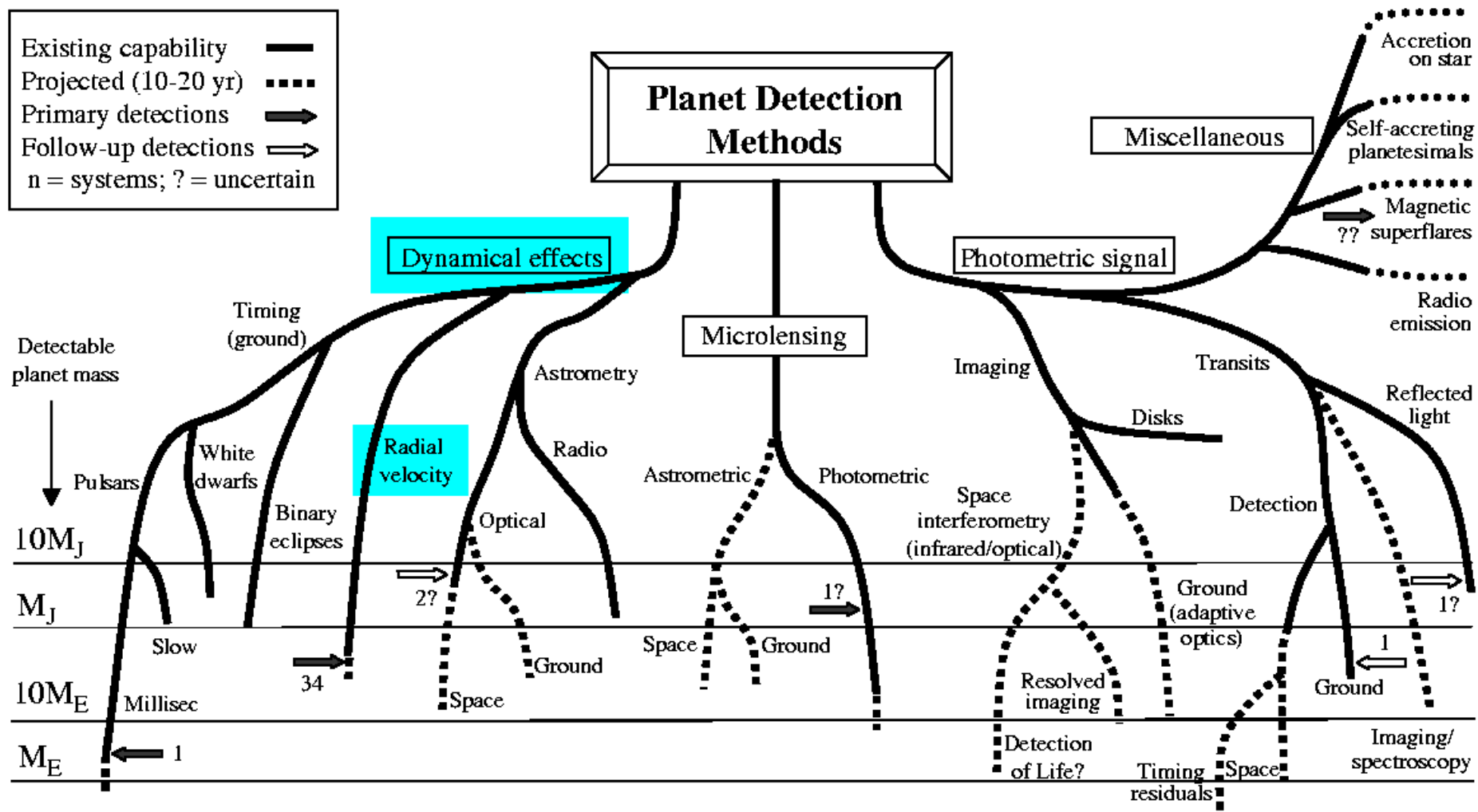
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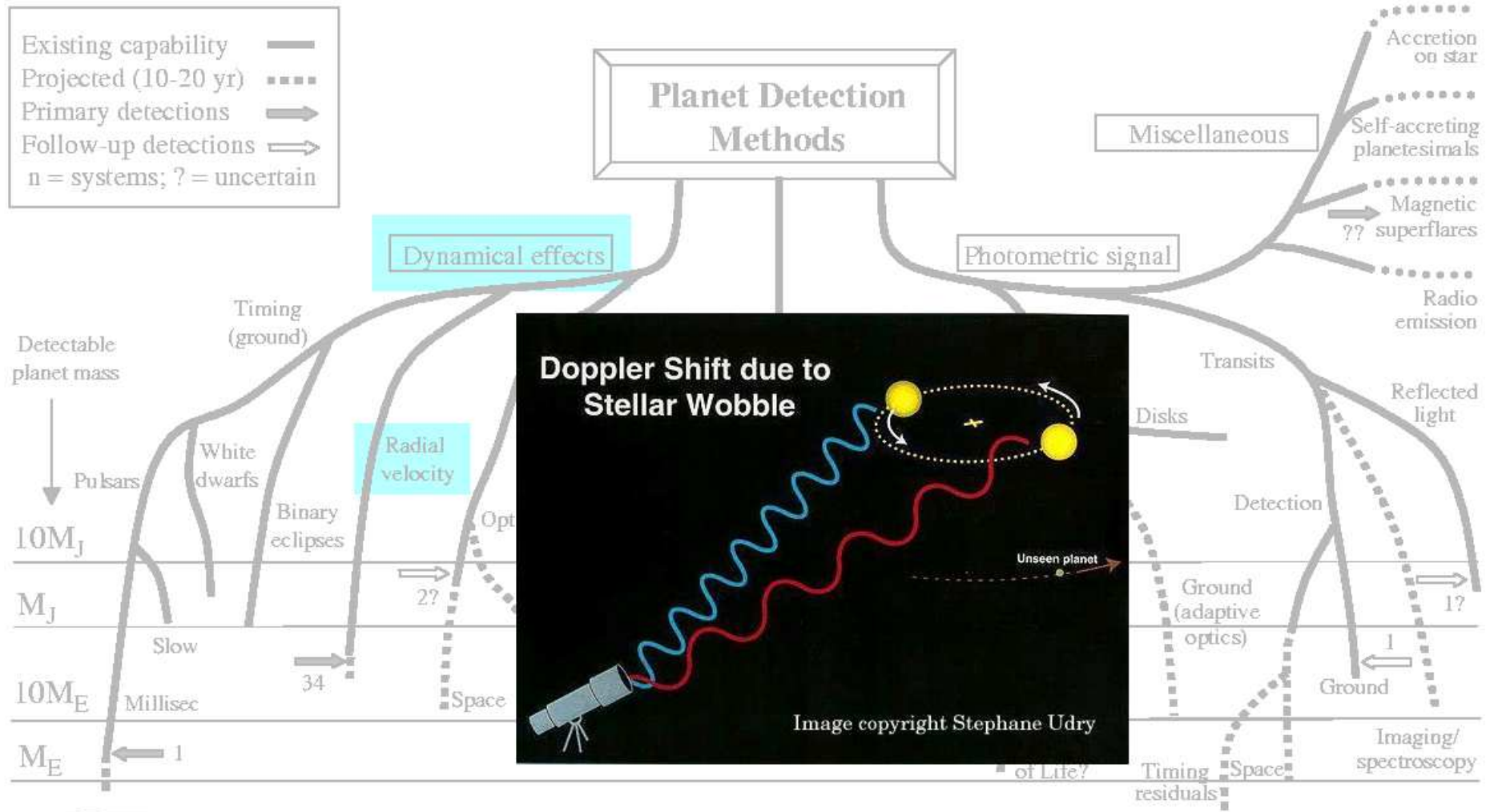


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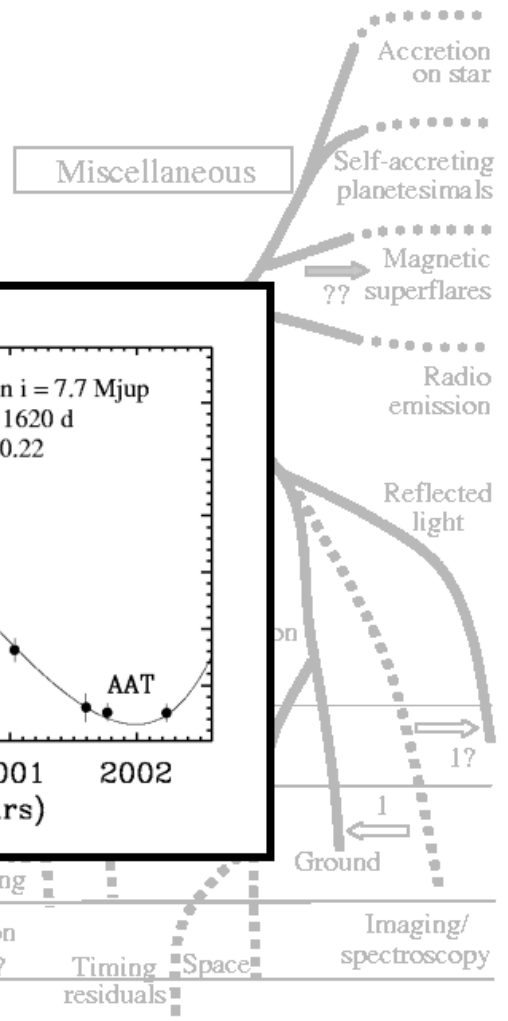
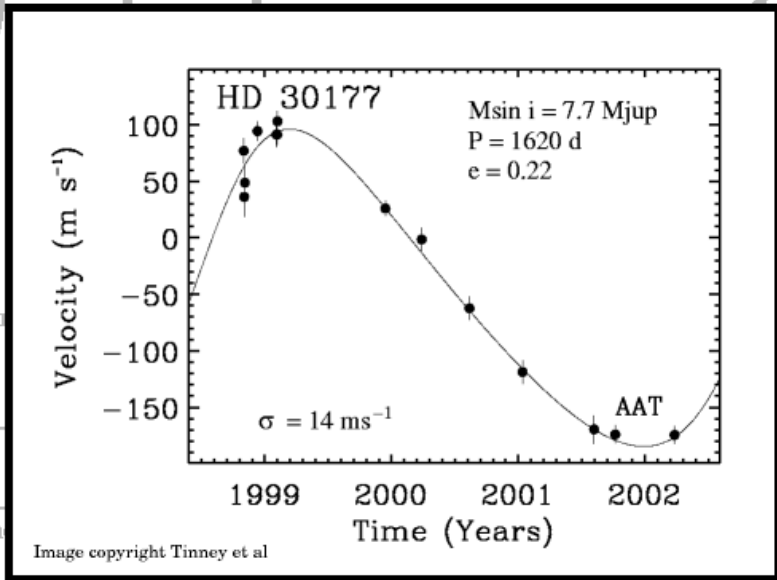
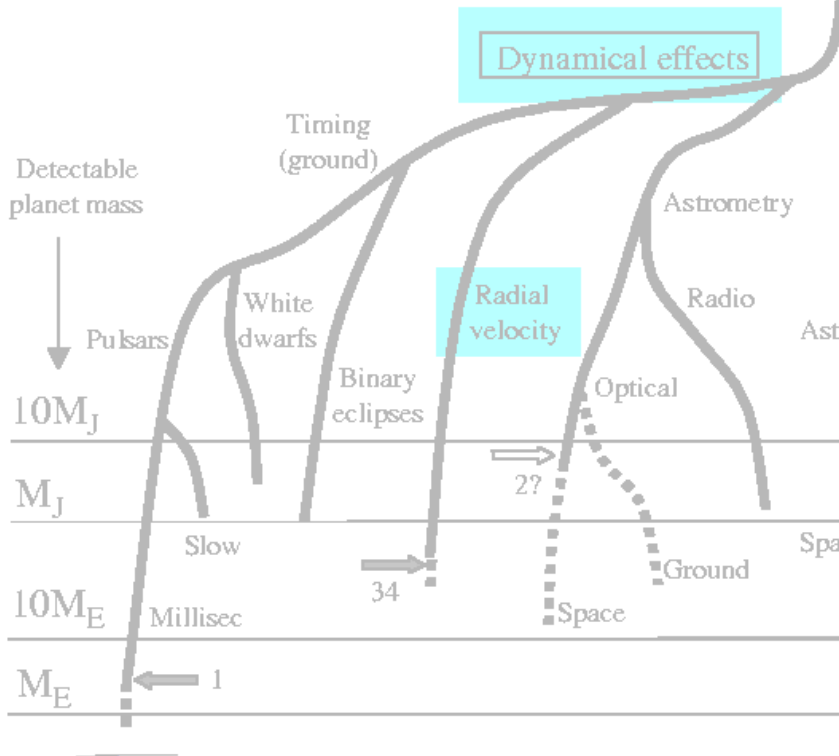
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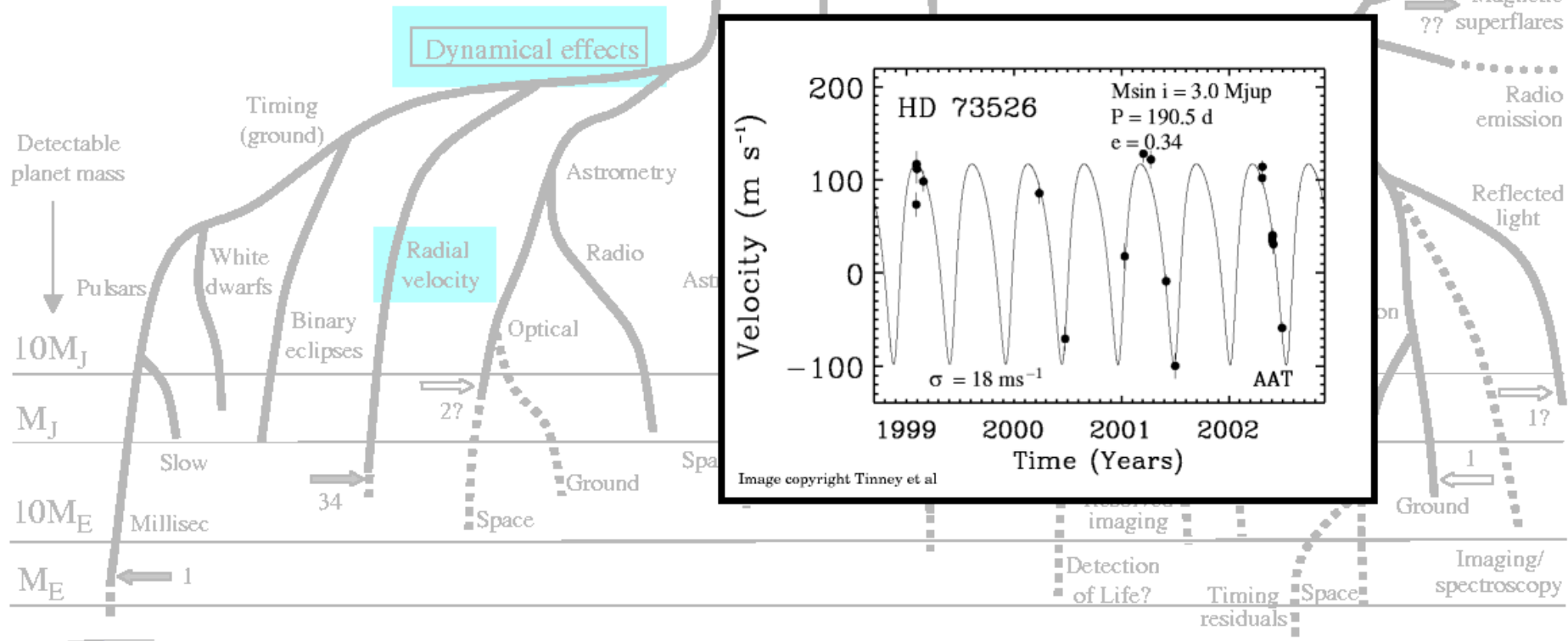
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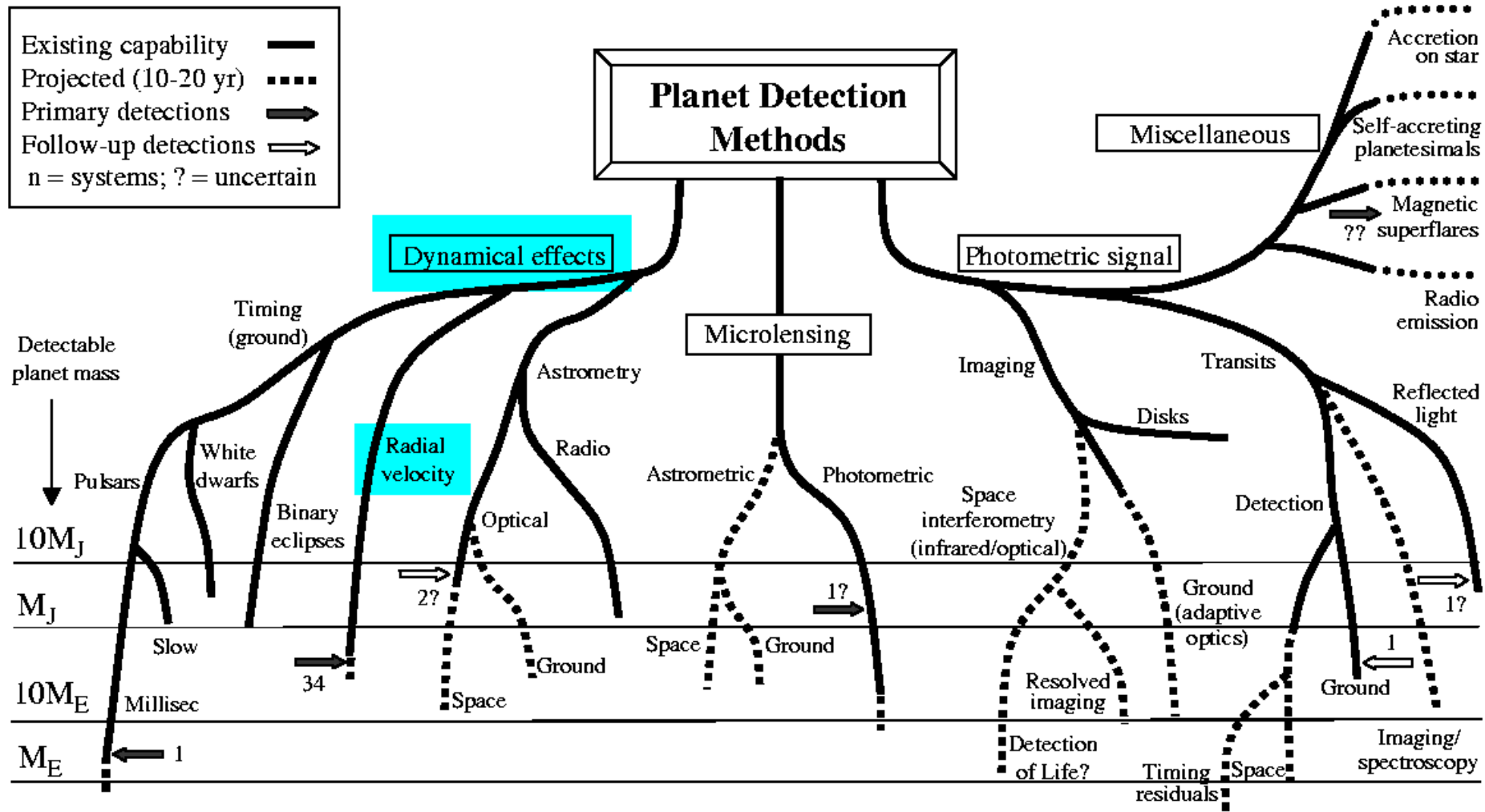
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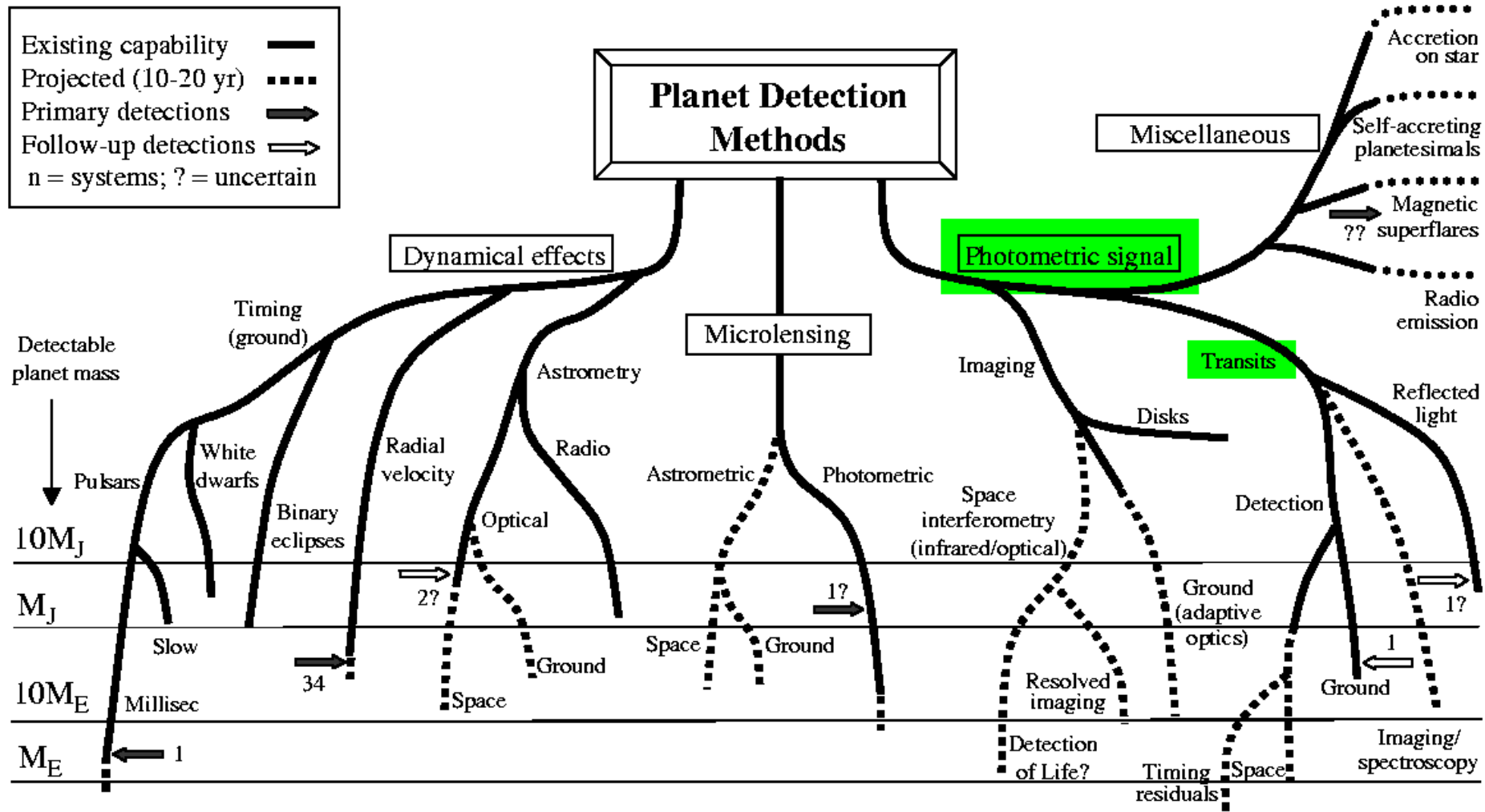
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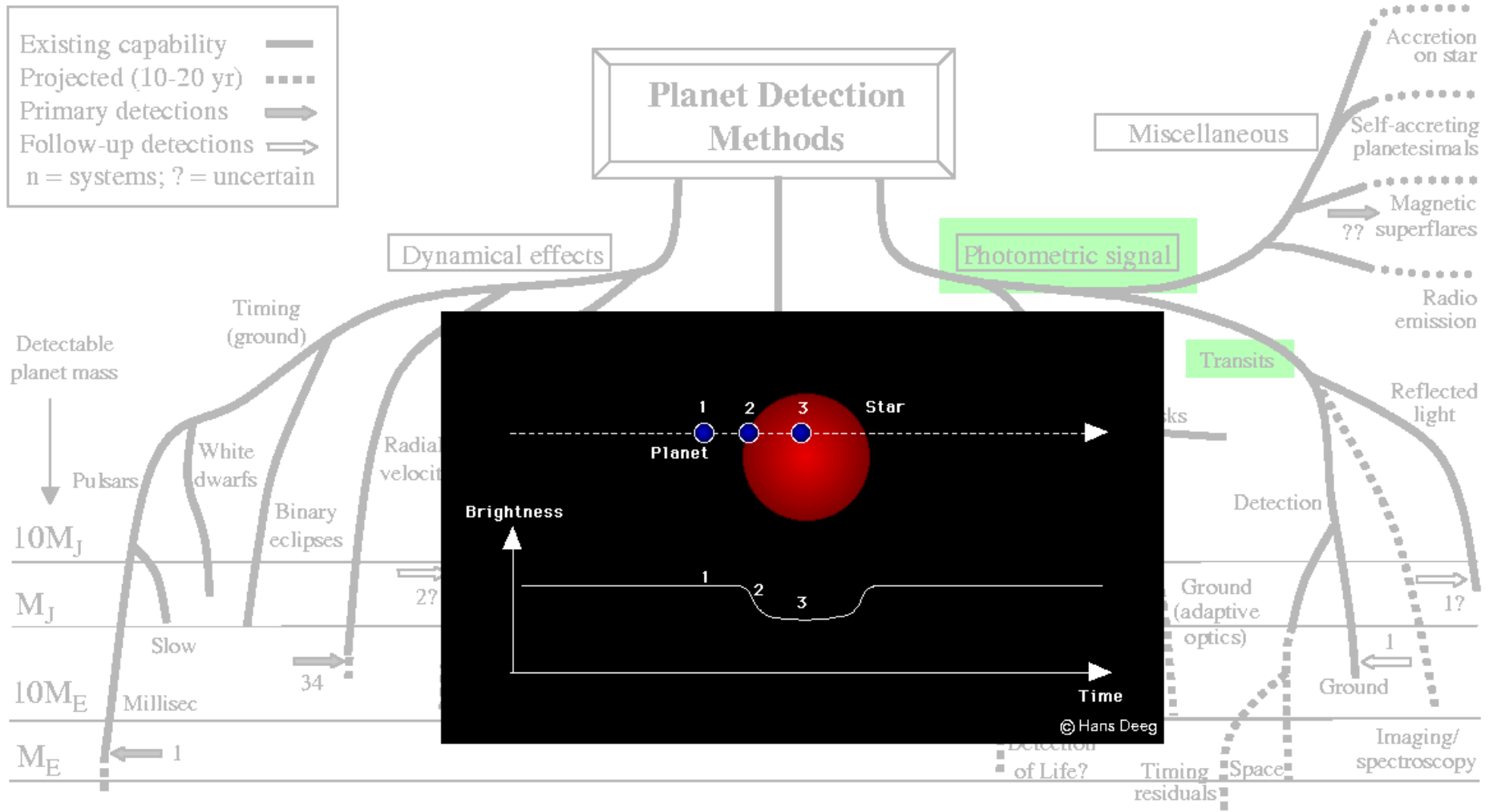
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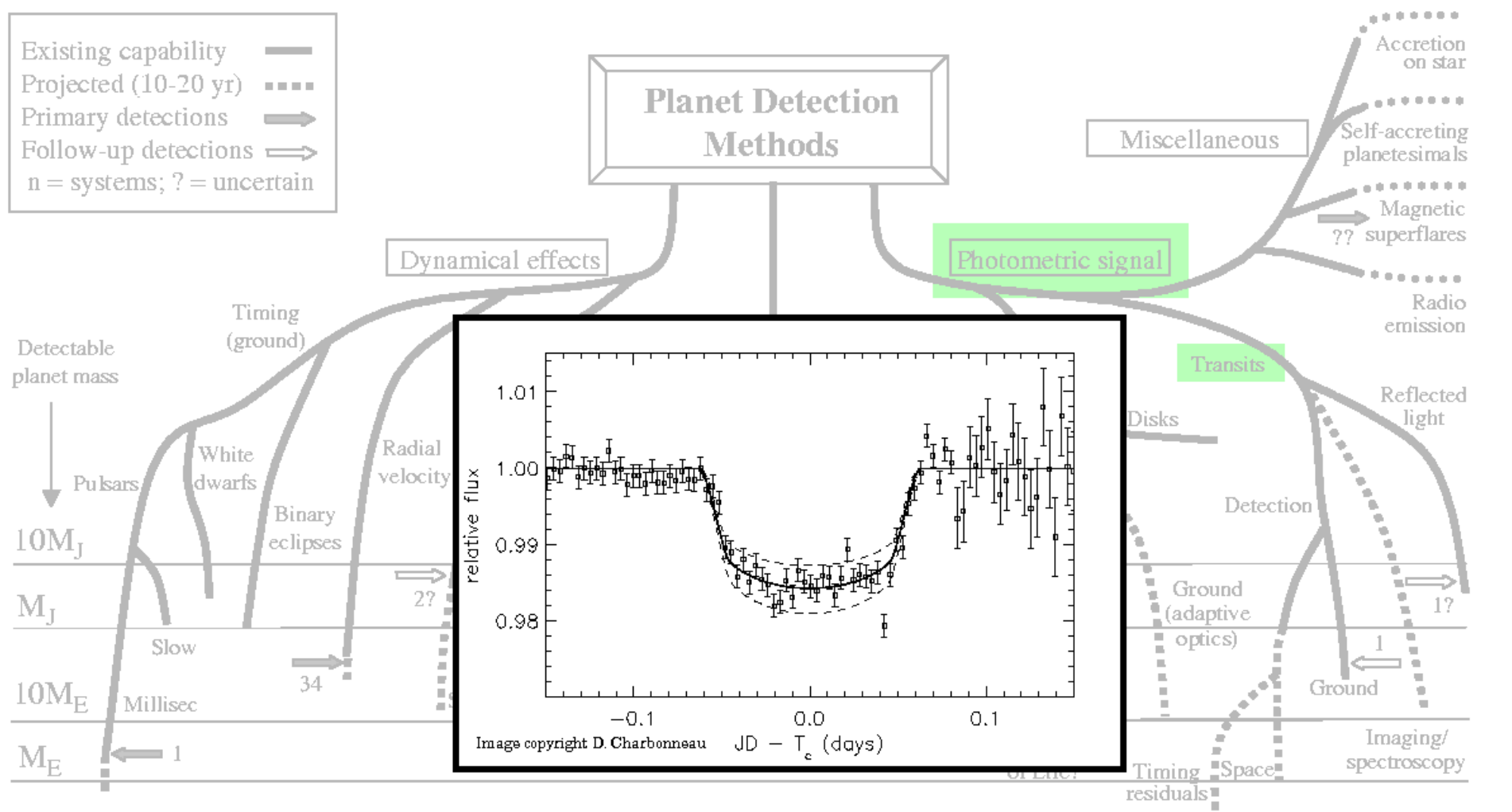


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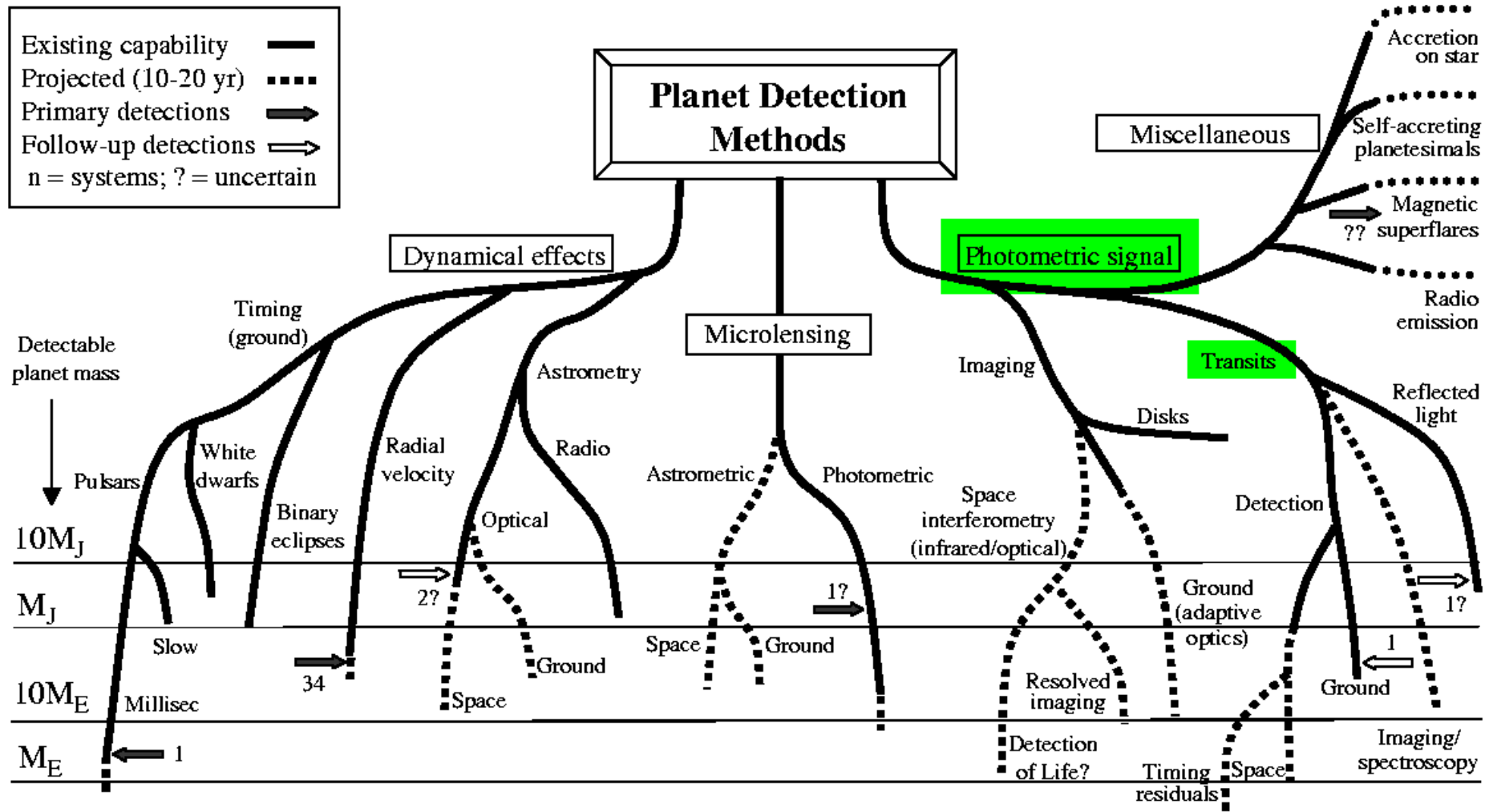


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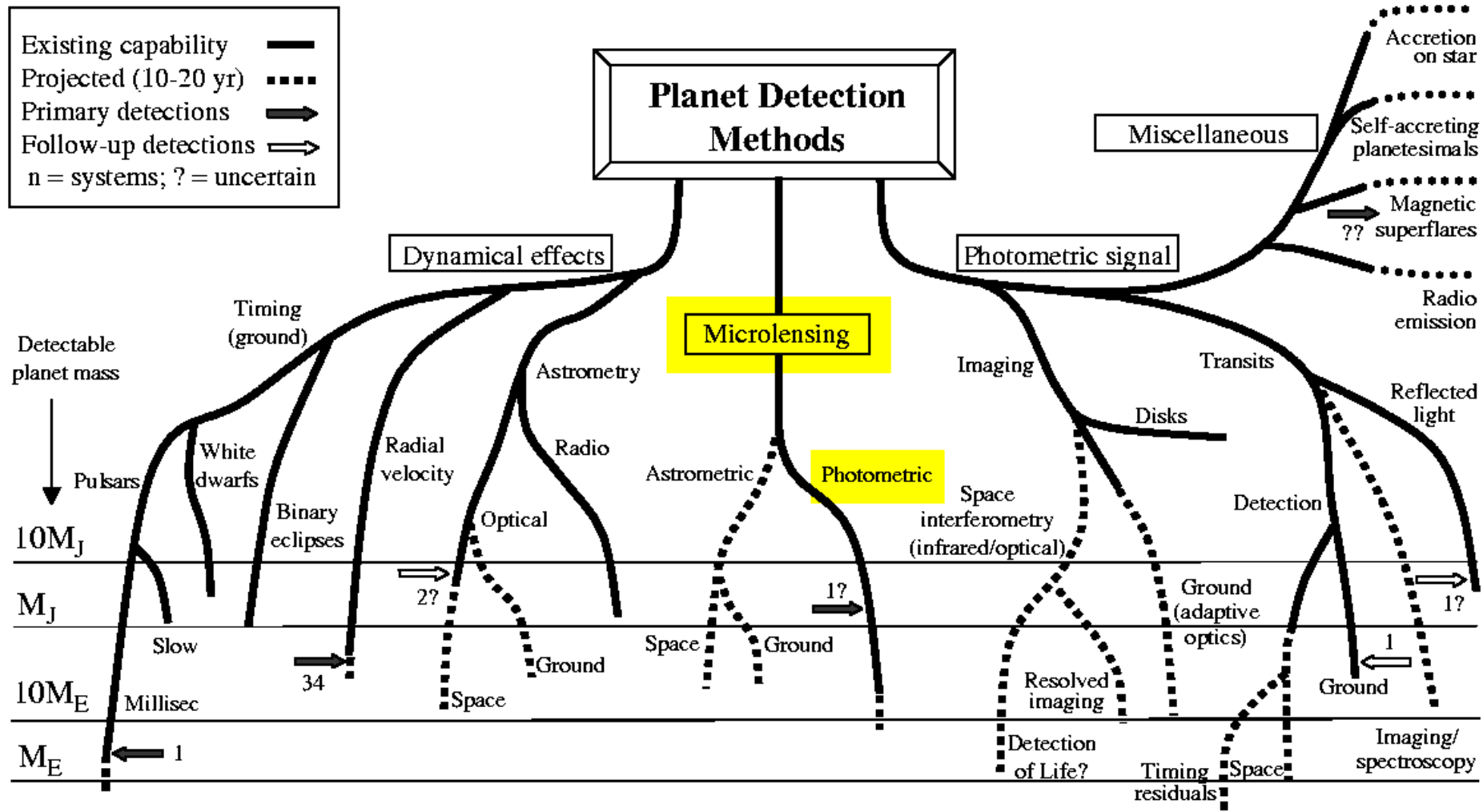
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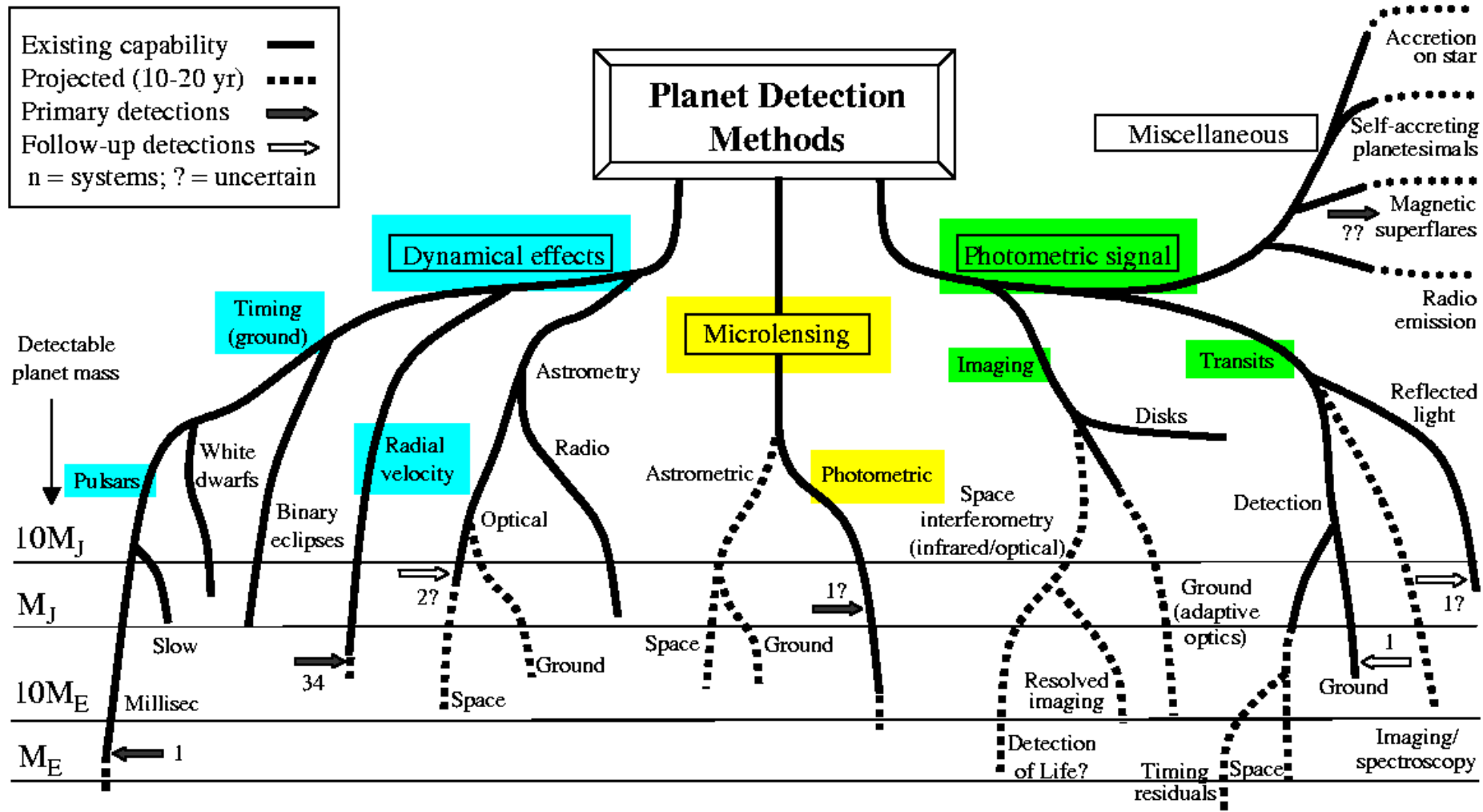
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Microlensing is currently detecting planets in a previously unreachable region of the planetary mass-radius space. Microlensing is returning detections of planets with masses approaching that of Earth.

Detection space

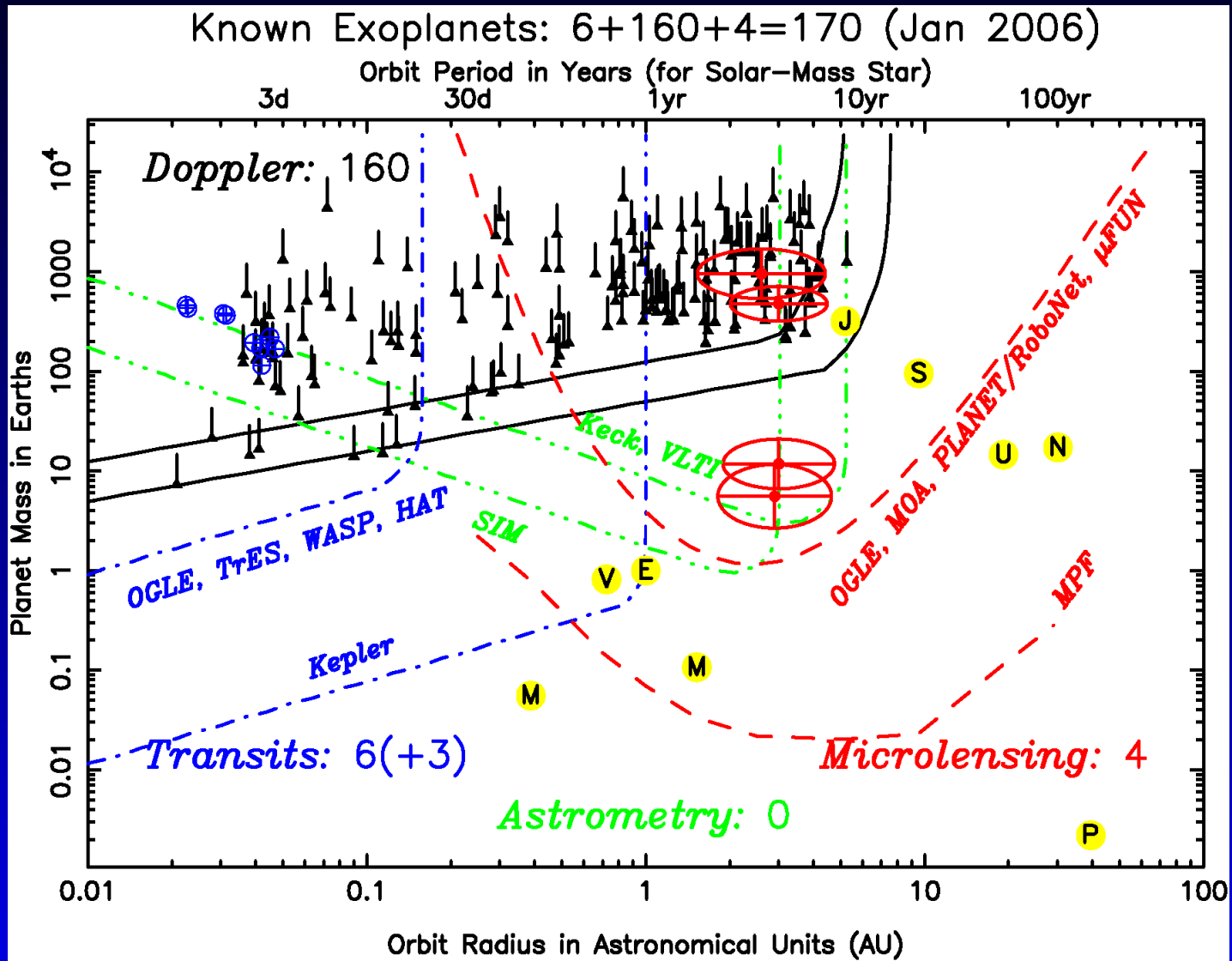


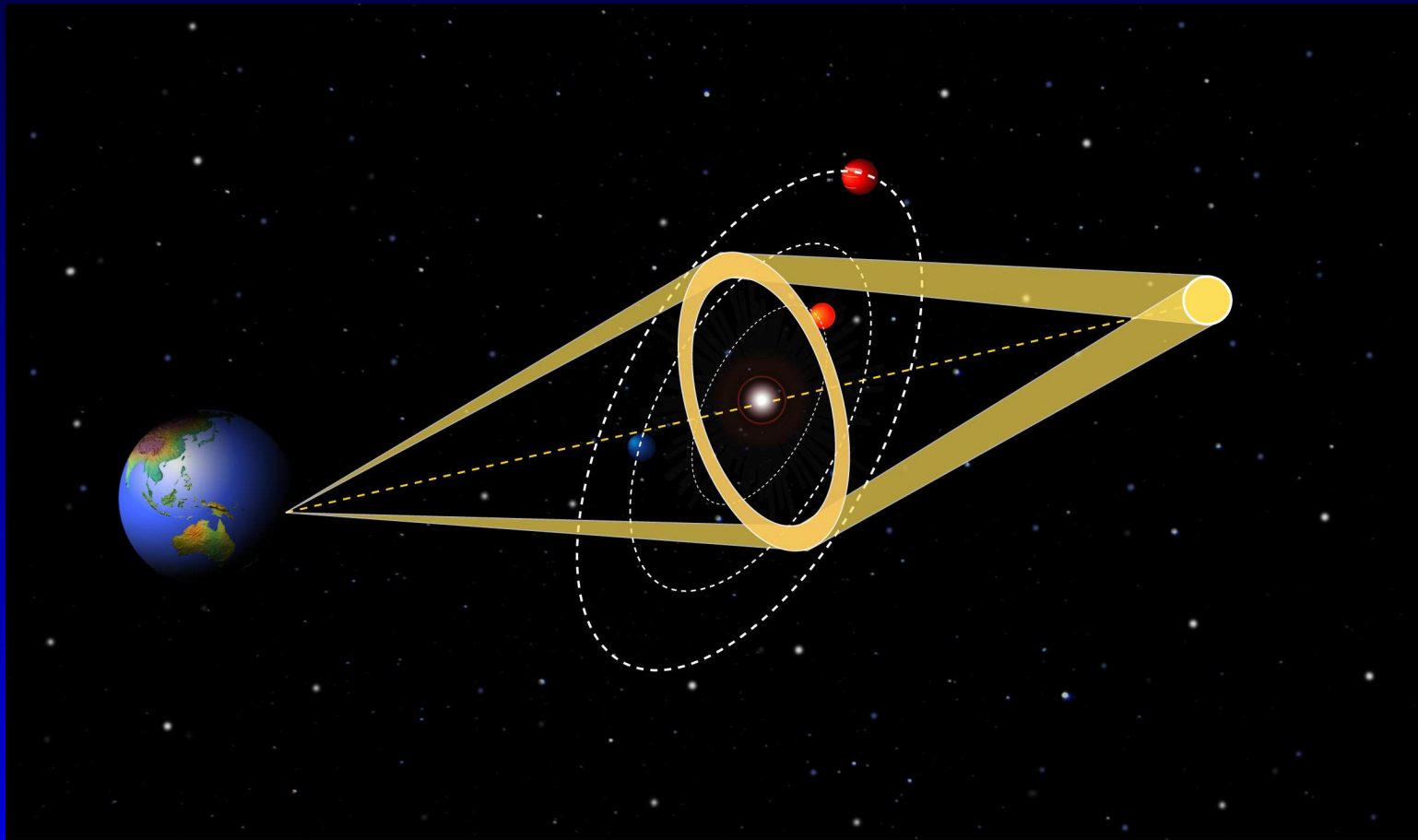
Figure courtesy **K. Horne**

Capability

The lensing effect is a “snapshot” of the lens system:
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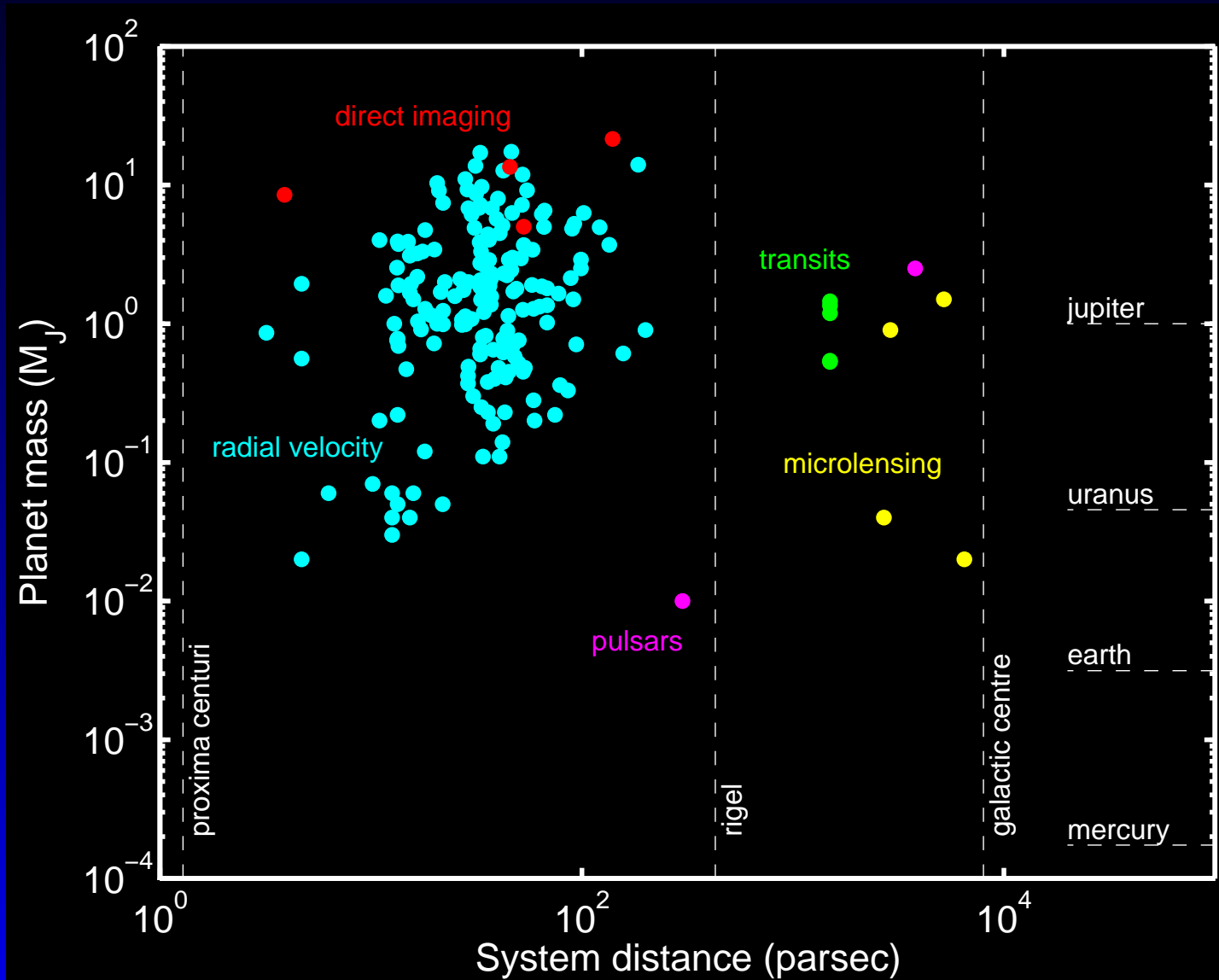
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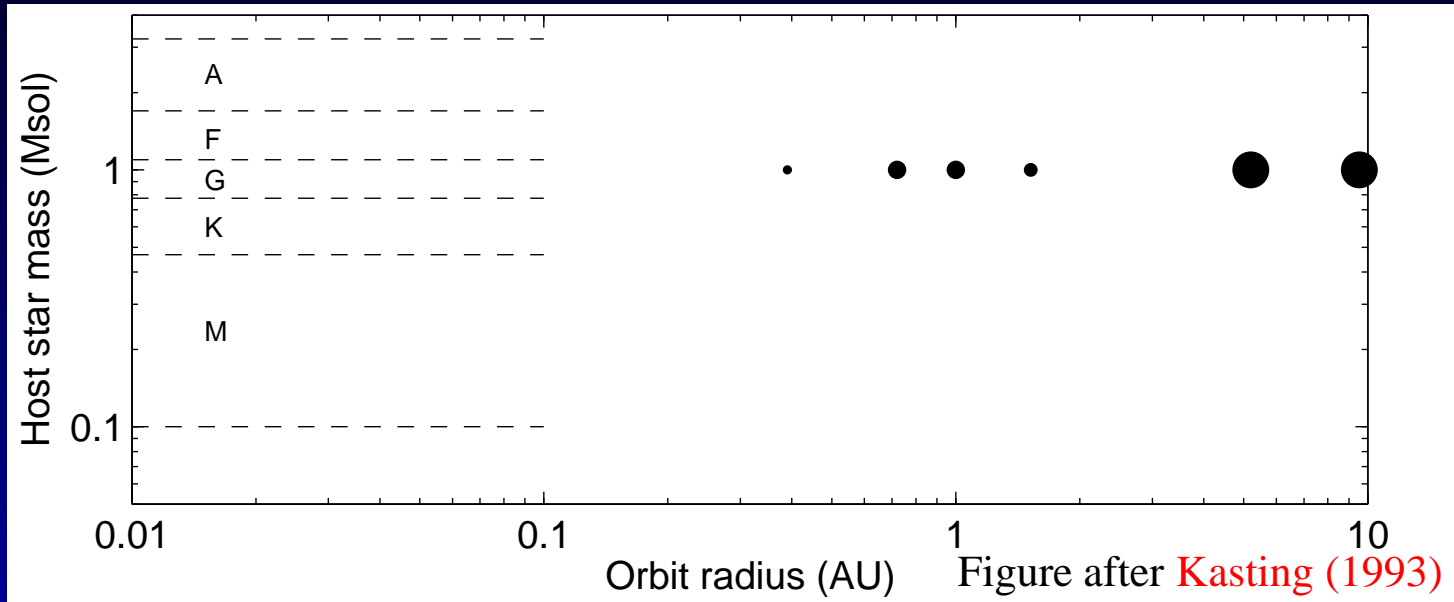
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Distant planetary systems are being discovered.

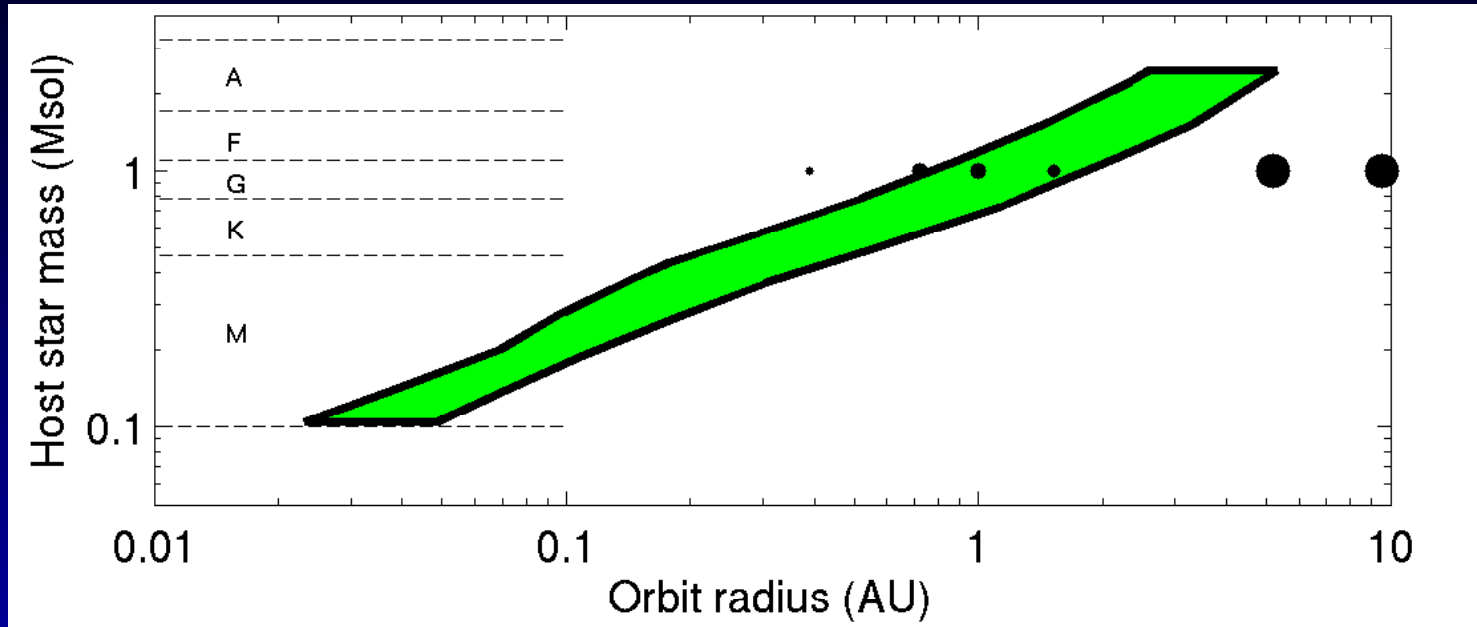
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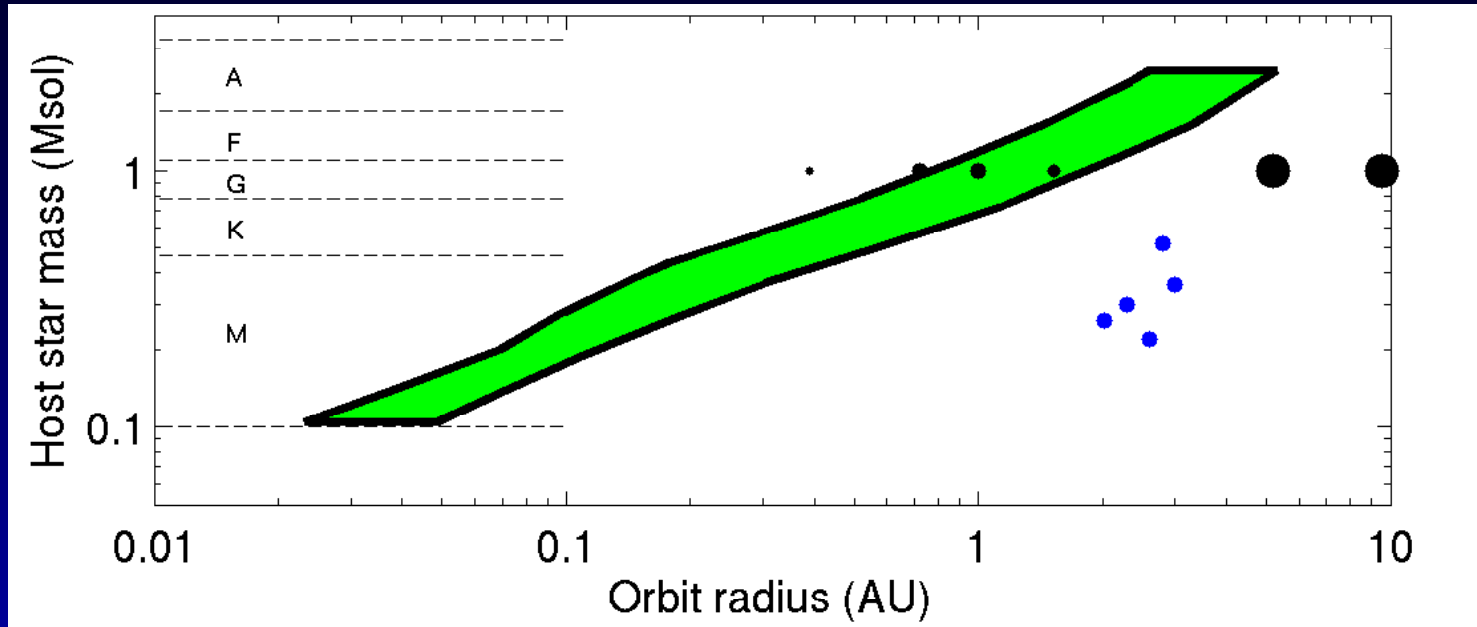
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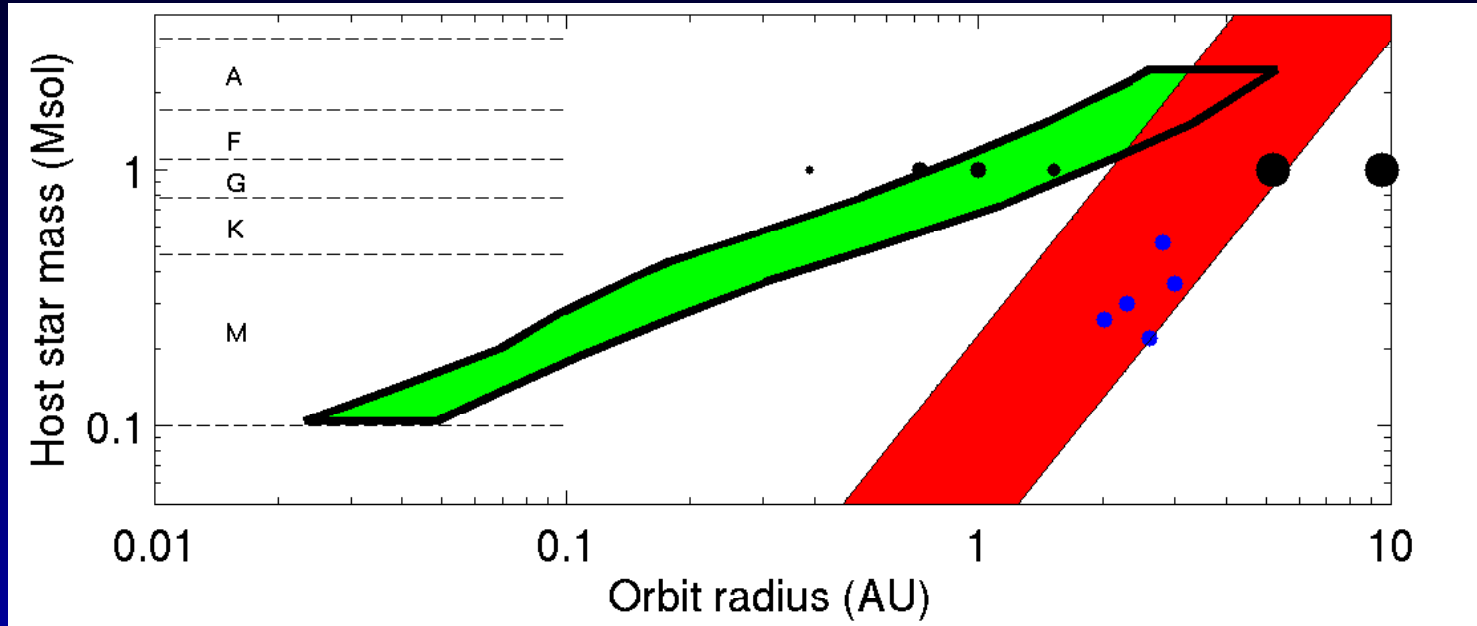
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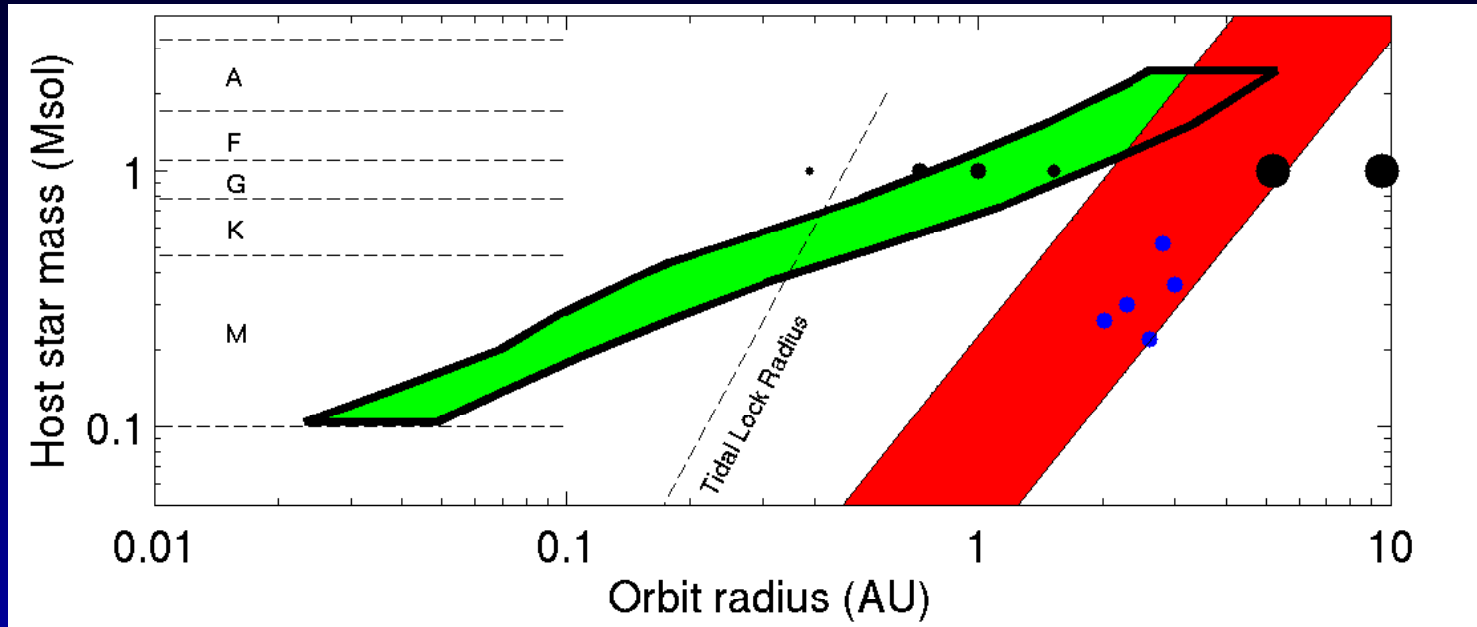
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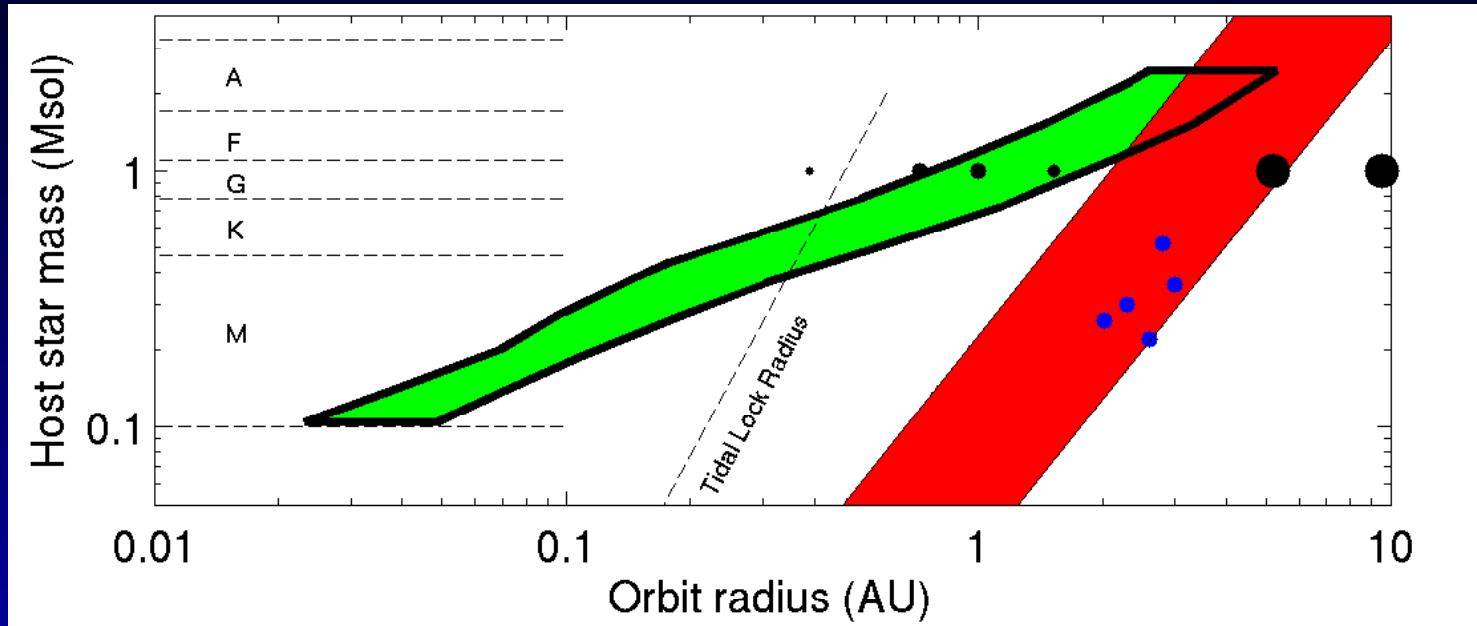
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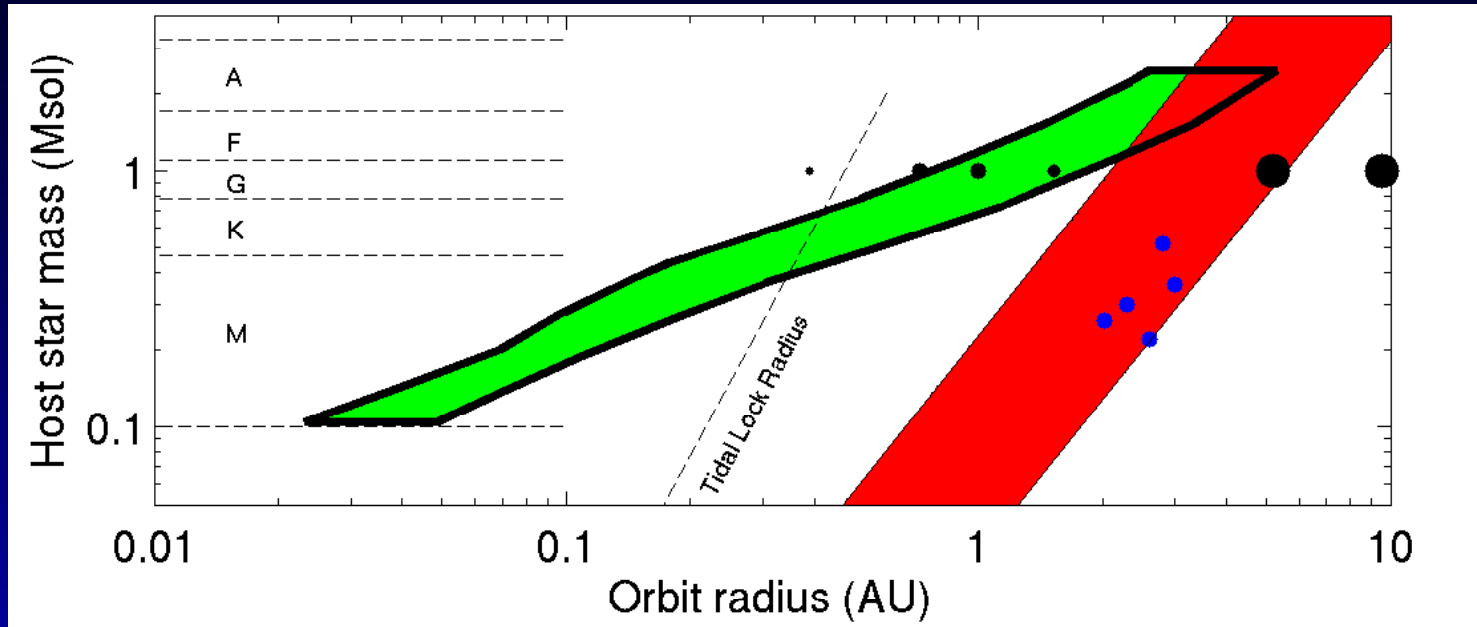


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However:

In studies of nearby solar type stars, *Abt & Levy (1976)* found 58% of the G dwarf primary stars had one or more stellar companions(see also *Abt, 1987*). *Duquennoy & Marcy (1991)* find a similar fraction of 57%. The multiplicity amongst M dwarfs is slightly less, at 42% (*Fischer & Marcy, 1992*).

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Most common “goodness-of-fit” parameter is

$$\chi^2 = \sum_{i=1}^N \frac{(y_i(t) - \hat{y}_i(t))^2}{\sigma_i^2}$$

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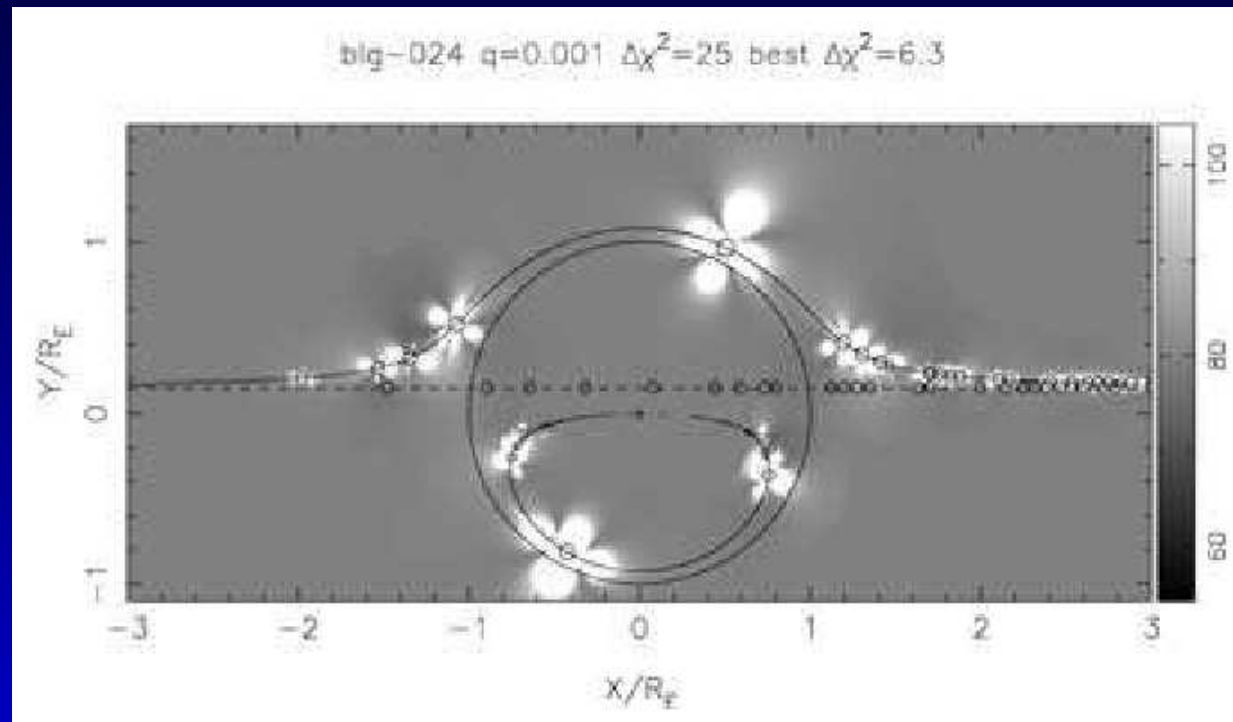
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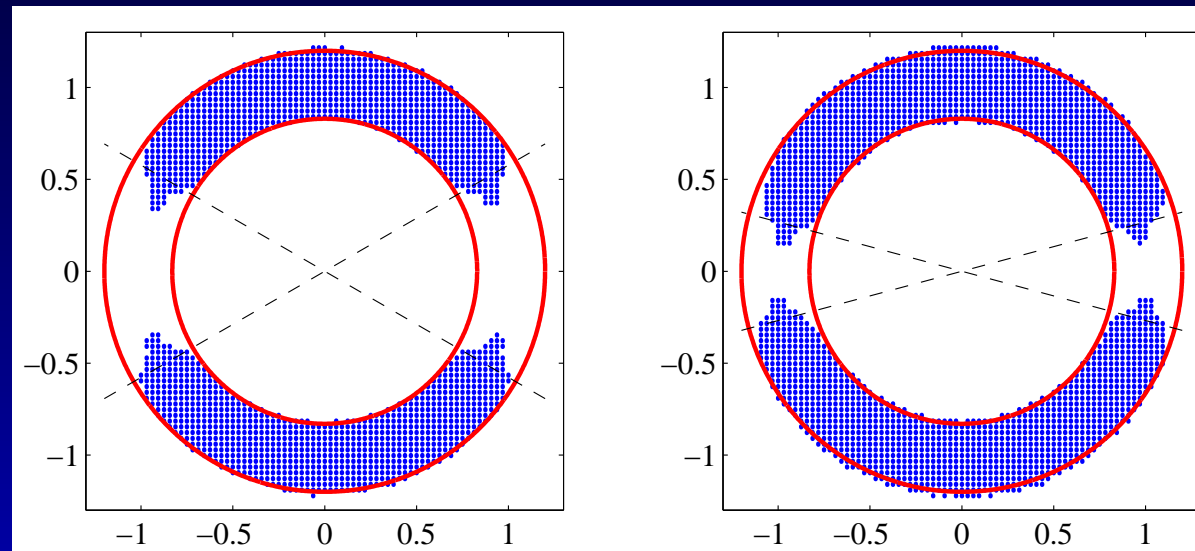


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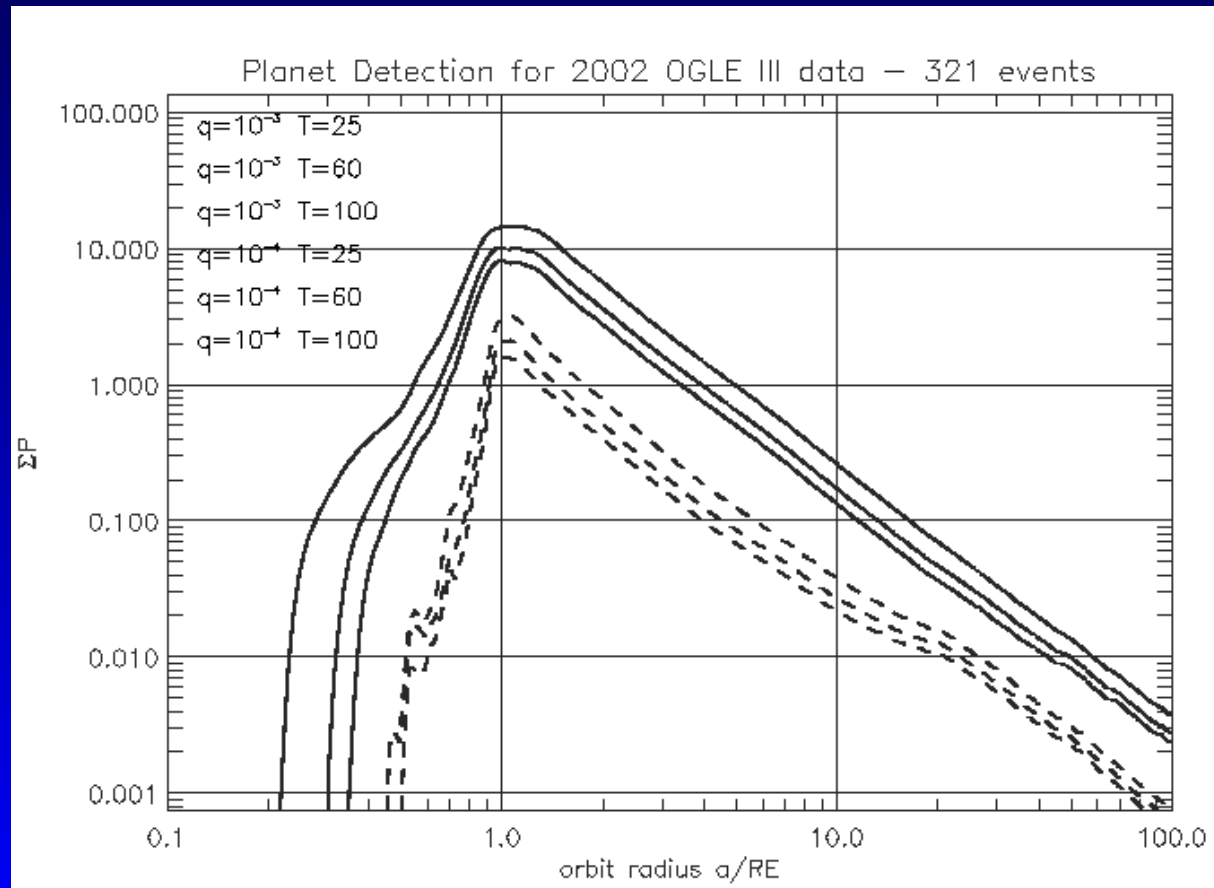
Detection limit maps for an Earth mass planet orbiting a $0.3M_{\odot}$ lens star. The axes are in units of R_E . The I band magnitude at maximum is 15, and the maximum amplification is 100. The χ^2 map on the left was generated using observation light curves comprised of 301 points over the interval $[-\frac{1}{2}t_{\text{FWHM}}, \frac{1}{2}t_{\text{FWHM}}]$. The right map was created using light curves comprised of 601 points over the interval $[-t_{\text{FWHM}}, t_{\text{FWHM}}]$.

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- Point-by-point significance - e.g 3 consecutive points deviating by $\geq 3\sigma$
- Coherency - deviating points follow a clear trend
- Confirmable - are deviating points supported by more observations

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But before we can apply detection criteria, we need to be able to fit microlensing lightcurve models!

Modelling

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The position of a planet is subject to a further degeneracy: models with projected planet distance a_p are degenerate with those with distance $1/a_p$.

Modelling

Point source single lens events are relatively easy to model. Standard non-linear fitting algorithms can be used to find u_{\min} , t_E and t_0 . The blend flux parameters F_u , F_l can similarly be found, either as additional parameters in the non-linear fitting routine, or through linear least-squares.

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$$u(t) = \left[u_{\min}^2 + \left(\frac{(t - t_0)}{t_E} \right)^2 \right]^{\frac{1}{2}}$$

$$\mu = \frac{u^2 + 2}{u\sqrt{u^2 + 4}}$$

$$F = F_l \cdot A(u(t)) + F_u$$

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Inverse ray shooting is a numerical technique which, although slow, is robust:

- (Rattenbury, 2002)

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The lightcurve generation code used depends on the required accuracy, speed and application.

The fitting algorithm used also needs to be well understood.

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The first coarse search could be a grid-wise search for example, or a random (uniform) sampling over the parameter space.

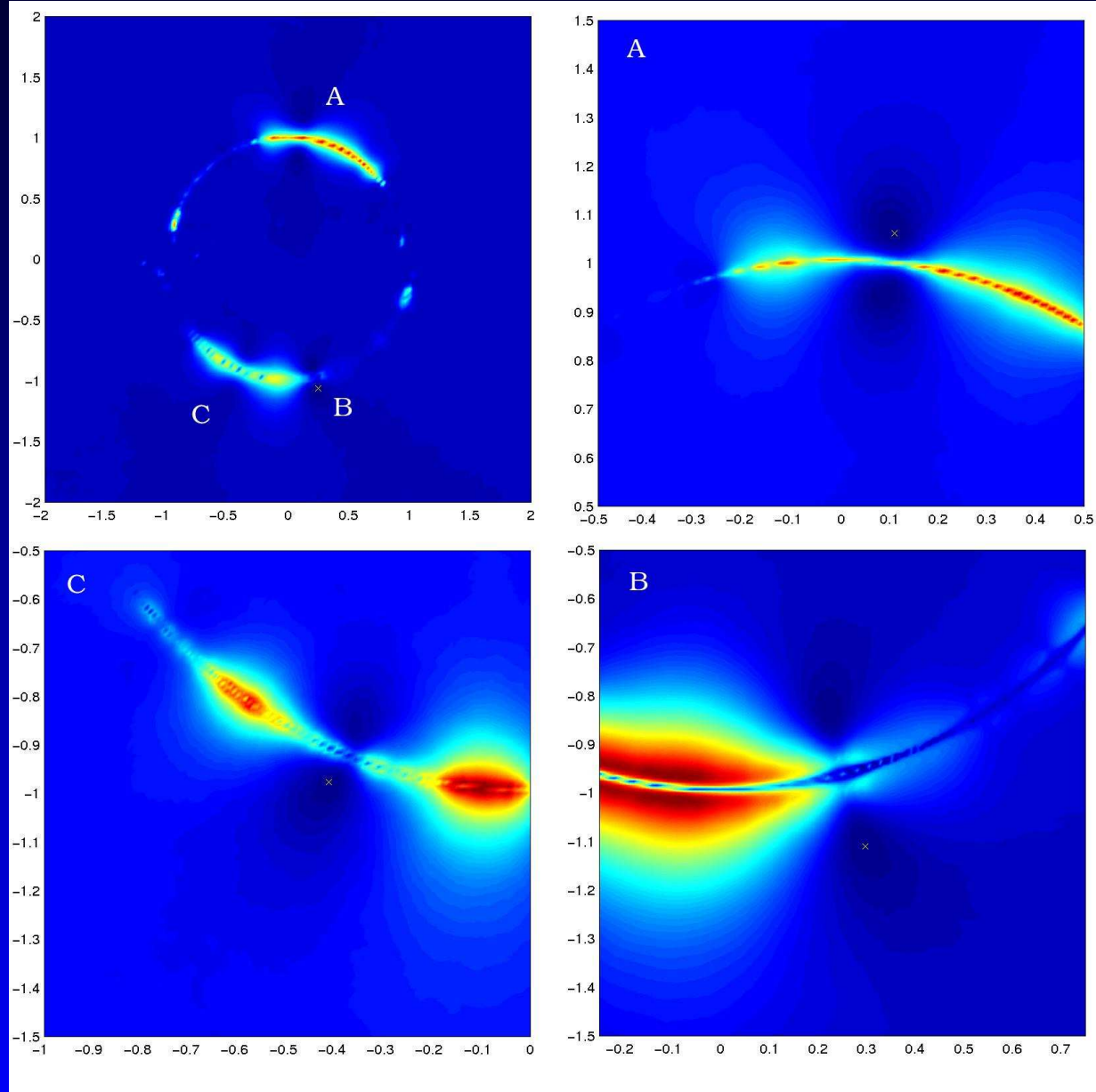
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May need several restarts, but should provide some good starting locations.

Modelling



Modelling

Having determined where possible optimal parameter values might lie in the parameter space

$\Omega = \Omega\{u_{\min}, t_0, t_E, r_s, x_{p_i}, y_{p_i}, q_i\}$, the next step is to apply more rigorous χ^2 minimisation procedures.

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Simple downhill simplex non-linear algorithms (Neadler-Mead) are subject to getting trapped in local χ^2 minima.

The simplex method always moving toward a lower χ^2 value. If the χ^2 manifold over the N-dimensional parameter space is not smooth, methods such as the downhill simplex can get trapped in non-optimal parameter space states.

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It is for this reason that it is suggested that the simplex method is restarted often, around the point in Ω considered to be the set of parameter values that minimises χ^2 . If the procedure recovers the same point after starting from a number of different points, the set of optimal parameters is considered secure.

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If the χ^2 manifold over Ω is sufficiently complicated, there is a greater risk that the results returned from the simplex methods are not global minima.

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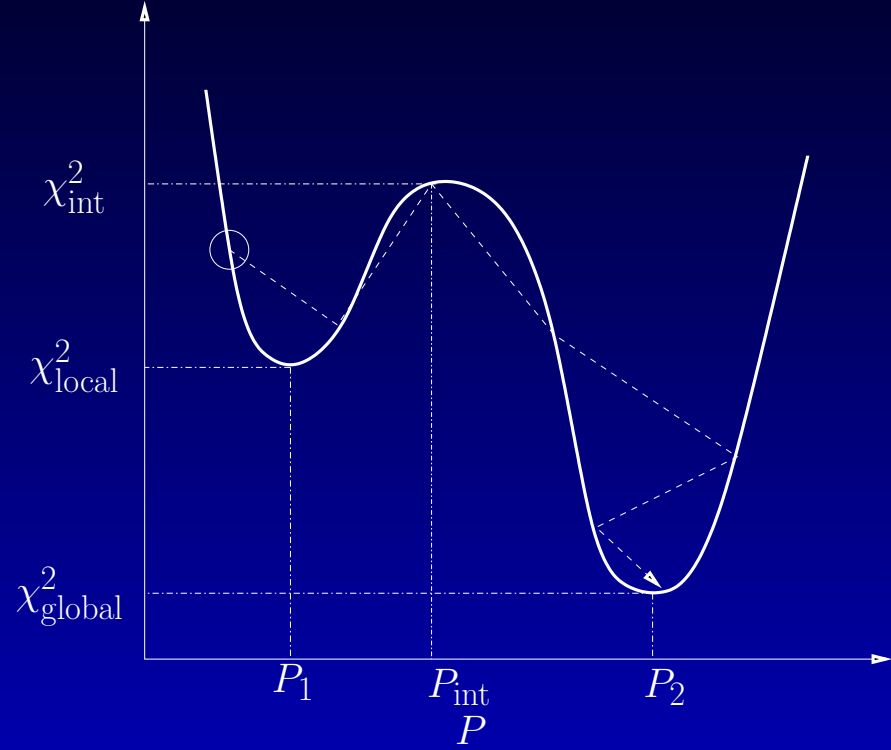
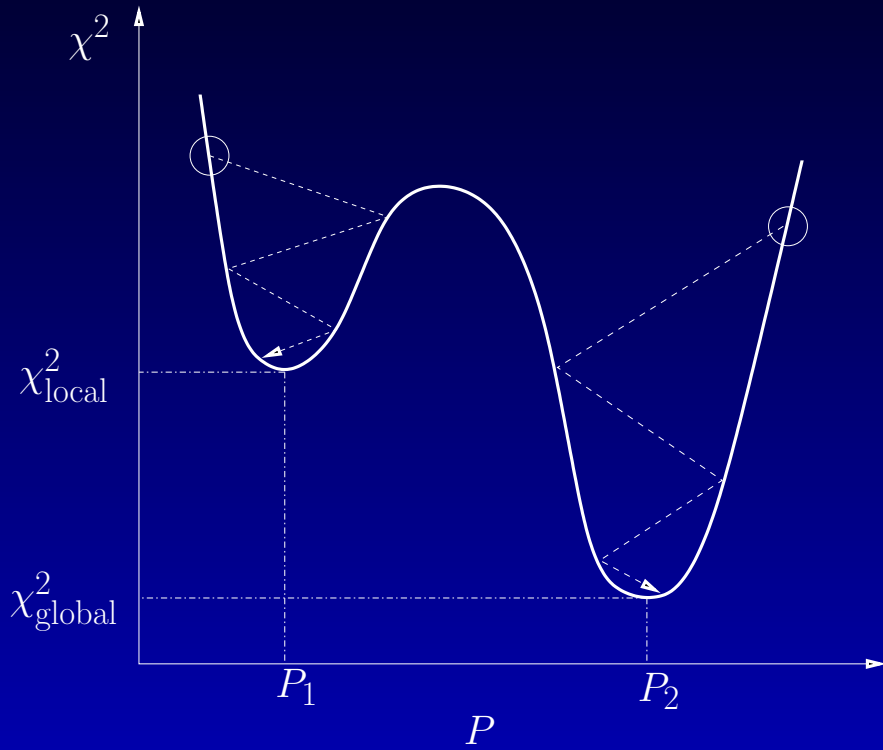
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This feature of the MHMCMC algorithm allows the method to occasionally disregard a (possibly better) candidate state. A chance therefore exists for the algorithm to find a more favourable region on the χ^2 manifold.

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Let $\{Z_k\}_{k=1}^M$ define a Markov chain, length M . Each element of Z_k is a set of parameters: $Z_k = \Omega_k$.

Modelling

The next element in the chain, Z_{k+1} is determined as follows:

1. Choose a candidate set of parameters, Ω' , by varying one or more of the parameters:

$$\begin{aligned}u_{\min}' &= u_{\min} + \mathcal{R}_{u_{\min}} \\t_0' &= t_0 + \mathcal{R}_{t_0} \\t_E' &= t_E + \mathcal{R}_{t_E} \\\vdots &= \quad \quad \quad \vdots\end{aligned} \tag{0}$$

where \mathcal{R} is an appropriately scaled random number.

Modelling

2. The candidate state, Ω' , is accepted with probability:

$$\alpha(\Omega'|\Omega) = \min \{1, \mathcal{F}(f'(t), f(t))\}$$

where $f(t)$ is the model function calculated using the current parameter set, Ω . $f'(t)$ is the model function, calculated using the candidate parameter set, Ω' . \mathcal{F} is a function based on the difference between the current and candidate models. If the candidate state Ω' is not accepted, set $Z_{k+1} = \Omega_k$.

Modelling

$\alpha(\Omega'|\Omega)$ is the acceptance ratio. The function \mathcal{F} essentially compares the χ^2 values arising from fitting the light curve generated using the candidate parameters Ω' , to that of the current set.

$$\mathcal{F} = \exp \left[\frac{\chi^2 - \chi'^2}{2} \right]$$

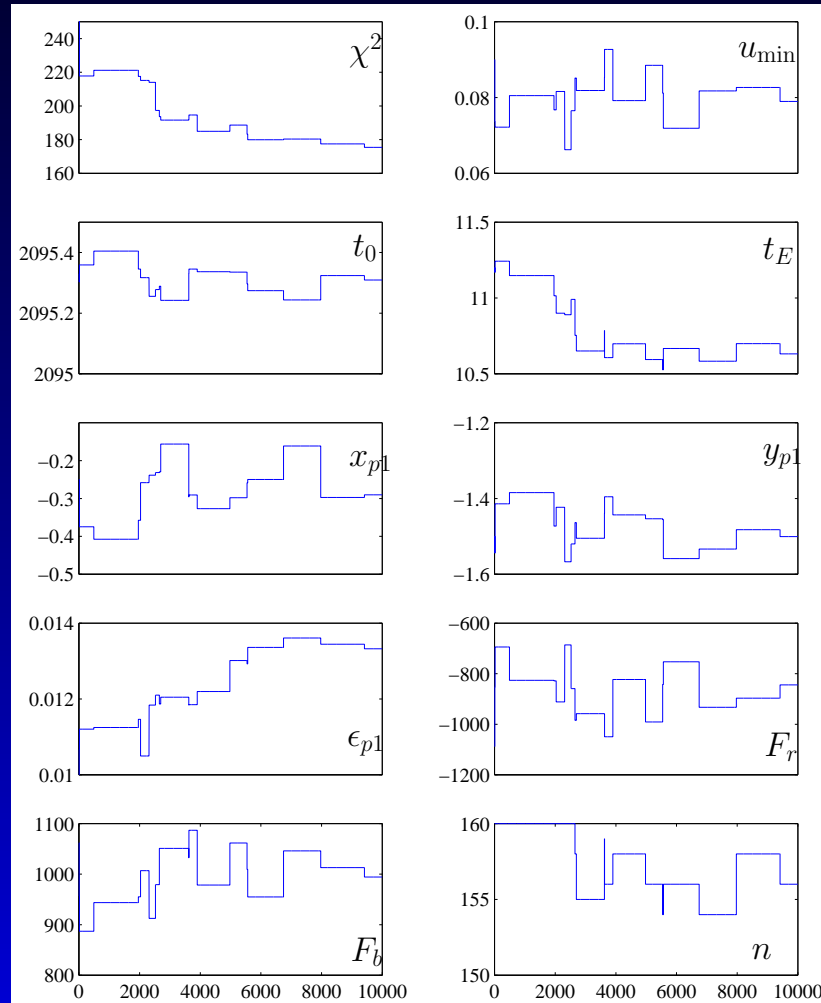
Modelling

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The above two procedures are iterated M times. Each parameter is allowed to vary, and the algorithm settles at a point in the parameter space which minimises χ^2 .

Modelling



Modelling

We can also impose an annealing schedule, $T(n)$:

$$\mathcal{F} = \exp \left[\frac{\chi^2 - \chi'^2}{2T} \right]$$

If $T \rightarrow 0$, the distribution will converge on the optimal modal space in the state space. Given a distribution in equilibrium, the “temperature”, T , can slowly be lowered and allow the distribution to find its new equilibrium. As $T \rightarrow 0$ the equilibrium will settle on the mode state(s).

Modelling

T	$\Delta\chi^2 > 0$		$\Delta\chi^2 < 0$	
	$\alpha(z' z)$	Accept z'	$\alpha(z' z)$	Accept z'
> 1	1	always	≤ 1	often
1.0	1	always	$0 < \alpha < 1$	sometimes
$0 < T < 1$	1	always	$0 < \alpha \ll 1$	seldom
0.0	1	always	0	never

$\Delta\chi^2 = \chi^2 - \chi'^2$, where χ'^2 corresponds to the candidate state, and χ^2 to the current state.

Modelling

- MHMCMC is excellent for fitting microlensing lightcurves
- MHMCMC fitting can be optimised in many ways.
- MHMCMC can be parallized easily.
- There is a lot of information in the equilibrium state of each parameter.

Genetic algorithms such as Particle Swarm Optimisation, Artificial Neural Networks and Self Organising Maps are other possible categorization methods.